

Matrix element of the electromagnetic operator between Kaon and pion states

Baum, Itzhak¹ Lubicz, Vittorio² Martinelli, Guido³
Simula, Silvano²

¹Rome University “La Sapienza”

²University of Rome III and INFN - Roma Tre

³University of Rome “La Sapienza” and INFN Rome

Villasimius, 15th June, 2010

Outline

Pion and
Kaon EM
matrix
elements

Baum et al.

Introduction

Calculations

Results

Work in
progress

Conclusions

1 Introduction - Kaon rare decays

2 Calculation of the EM operator matrix element

3 Results

4 Work in progress

5 Conclusions

Kaon rare semileptonic decays as new physics probes

Pion and
Kaon EM
matrix
elements

Baum et al.

Introduction

Calculations

Results

Work in
progress

Conclusions

Rare Kaon decays have not been detected yet:

$$BR(K_L \rightarrow \pi^0 \ell^\pm \ell^\mp)_{exp} < 6.6 \cdot 10^{-10}$$

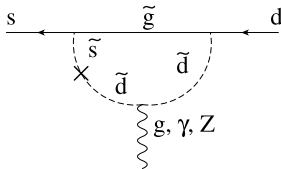
In the SM they are estimated to be

$$BR(K \rightarrow \pi e e)_{SM} \sim 1.5 \cdot 10^{-12}$$

$$BR(K \rightarrow \pi \mu \mu)_{SM} \sim 3 \cdot 10^{-10}$$

New physics can be the leading contribution, mediated through the Electro-magnetic and Chromo-magnetic operators:

$$\begin{aligned} \mathcal{Q}_{EM}^+ &= \bar{s} F_{\mu\nu} \sigma^{\mu\nu} d \\ \mathcal{Q}_{CM}^+ &= \bar{s} G_{\mu\nu} \sigma^{\mu\nu} d \end{aligned}$$



- Sensitive to hadronic matrix elements.

Previous work

Pion and
Kaon EM
matrix
elements

Baum et al.

Introduction

Calculations

Results

Work in
progress

Conclusions

First lattice calculation [Becirevic et al. 2001] of the EM form factor

$$f_T(q^2 = 0) = 0.77 \pm 0.06 \pm 0.03$$

With the slope in q^2

$$\lambda = 1.21 \pm 0.05 \text{ GeV}^{-2}$$

- Quenched ($n_f = 0$)
- High pion masses ($530 < m_\pi < 800 \text{ MeV}$)
- One lattice size ($a^{-1} = 2.7(1) \text{ GeV}$)

Lattice details

Pion and
Kaon EM
matrix
elements

Baum et al.

Introduction

Calculations

Results

Work in
progress

Conclusions

- ETMC lattice QCD simulations [ETMC 0701012, 0911.5061]
- Dynamical flavors: $n_f = 2$
- Pion mass range: $270 < m_\pi < 600$ MeV
- Lattice sizes: $24^3 \times 48$ and $32^3 \times 64$
- Lattice step sizes: $a = 0.068, 0.085, 0.10$ fm
- Action is Symanzik tree-level improved with maximally twisted-mass Wilson fermions
- Non perturbative renormalization in the RI/MOM scheme [ETMC 1004.1115]
- 3-point correlators with all-to-all stochastic propagator calculation, increase accuracy
- Breit momentum frame: $\vec{p}_K = \vec{p}, \vec{p}_\pi = -\vec{p}$

Electromagnetic form factor calculation

Pion and
Kaon EM
matrix
elements

Baum et al.

Introduction

Calculations

Results

Work in
progress

Conclusions

$$\mathcal{Q}_{EM} = \bar{s}\sigma^{\mu\nu}d$$

The EM form factor is acquired from the EM matrix element by [Becirevic et al. 2001]:

$$\left\langle \frac{\pi^0}{\sqrt{2}} | \mathcal{Q}_{EM} | K^0 \right\rangle = i (p_K^\mu p_\pi^\nu - p_K^\nu p_\pi^\mu) \frac{\sqrt{2}f_T}{m_K + m_\pi}$$

To obtain the matrix elements from the 3-point correlators, we look at the lattice times far from the pion and Kaon sources

$$C_3^{K\pi} \rightarrow \frac{\sqrt{Z_K Z_\pi}}{4E_K E_\pi} \langle \pi^0 | \mathcal{Q}_{EM} | K^0 \rangle e^{-E_K t_x - E_\pi (t_y - t_x)}$$

and use the ratio

$$\frac{C_3^{K\pi} C_3^{\pi K}}{C_2^\pi(t_y) C_2^K(t_y)} \rightarrow \frac{\langle \pi^0 | \mathcal{Q}_{EM} | K^0 \rangle^2}{16E_K E_\pi}$$

where t_y is a fixed point $t_y = T/2$.

Calculations

Pion and
Kaon EM
matrix
elements

Baum et al.

Introduction

Calculations

Results

Work in
progress

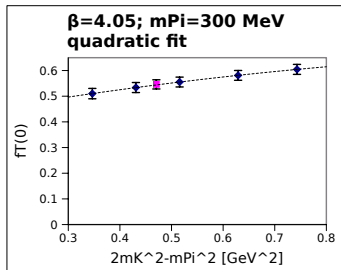
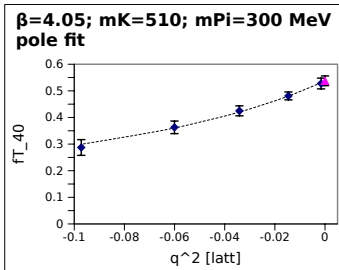
Conclusions

- Interpolation in momentum to $q^2 = 0$, assuming pole behaviour:

$$f_T(q^2) = \frac{f_T(0)}{1 - q^2\lambda}$$

- Interpolation to physical strange mass:

$$(2m_K^2 - m_\pi^2)_{LATT} \rightarrow (2m_K^2 - m_\pi^2)_{PHYS} \propto (m_s)_{PHYS}$$



Extrapolation in masses

Pion and
Kaon EM
matrix
elements

Baum et al.

Introduction

Calculations

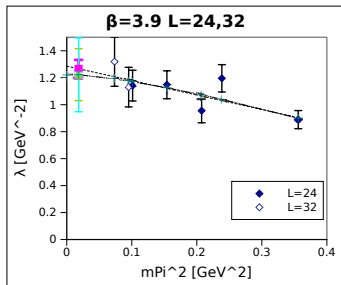
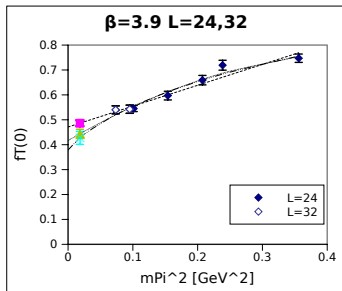
Results

Work in
progress

Conclusions

Extrapolation in m_π^2 to physical pion mass $m_{\pi^0} = 135$ MeV

- linear $f = Am_\pi^2 + B$
- quadratic $f = A'm_\pi^4 + B'm_\pi^2 + C'$
- log-linear $f = A''m_\pi^2 \ln(m_\pi^2) + B''m_\pi^2 + C''$



- Small finite volume effects

Results

Pion and
Kaon EM
matrix
elements

Baum et al.

Introduction

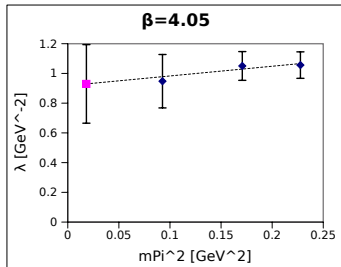
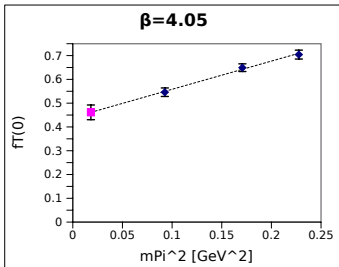
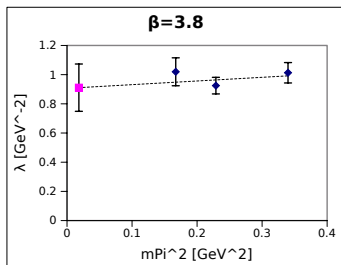
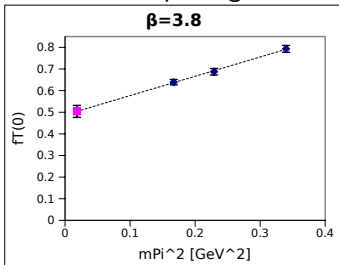
Calculations

Results

Work in
progress

Conclusions

Other lattice spacings:



Results

Pion and
Kaon EM
matrix
elements

Baum et al.

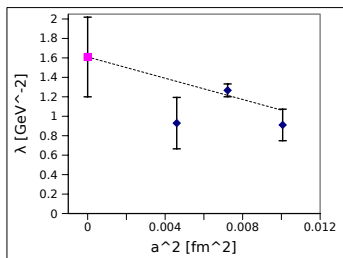
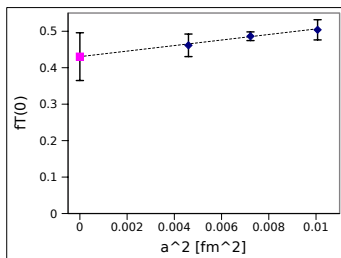
Introduction

Calculations

Results

Work in
progress

Conclusions



We find (Preliminary results, no systematic effects)

$$f_T(q^2 = 0) = 0.430 \pm 0.066^{stat}$$

$$\lambda = 1.61 \pm 0.41^{stat} \text{ GeV}^{-2} \text{ (slope in } q^2, \text{ pole fit)}$$

To compare with [Becirevic et al. 2001] (linear fit)

$$f_T(0) = 0.77 \pm 0.06 \pm 0.03$$

$$\lambda = 1.21 \pm 0.05 \text{ GeV}^{-2}$$

Results

Pion and
Kaon EM
matrix
elements

Baum et al.

Introduction

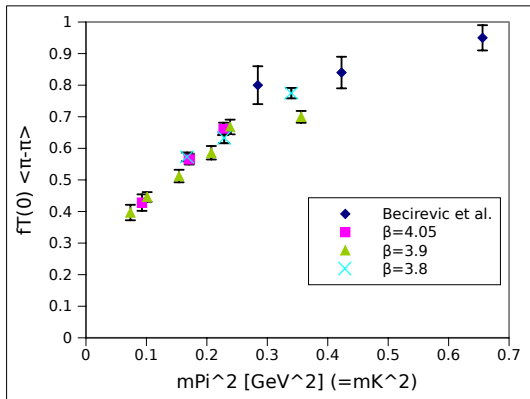
Calculations

Results

Work in
progress

Conclusions

Confronting results for $m_K = m_\pi$ for similar lattice sizes a



Similar behaviour, difference may be due to:

- Extrapolation from large pion masses
- Quenching effects

Work in progress

- Electro-magnetic operator $Q_{EM}^+ = \bar{s}\sigma^{\mu\nu}d$
 - Combined fit for all lattice spacings
 - Systematic errors analysis (chiral extrapolation, momentum dependence)
- Chromo-magnetic operator $Q_{CM}^+ = \bar{s}G_{\mu\nu}\sigma^{\mu\nu}d$
 - No previous lattice calculation
 - Matrix elements calculation
 - Renormalization

$$Q_{CM}^{renorm} = Z_{CM} \left(Q_{CM}^{bare} + \frac{c}{a} Q_S \right)$$

- additive – subtraction of mixing with scalar operator
- multiplicative – 1-loop lattice perturbation theory [H. Panagopoulos et al.]

Conclusions

Pion and
Kaon EM
matrix
elements

Baum et al.

Introduction

Calculations





Results

Work in
progress

Conclusions

- Previously, single calculation (2001) of the EM operator
- Our calculations were performed for a large range of masses and lattice spacings
- Higher statistical accuracy achieved
- Values at $q^2 = 0$ differ, may be due to either quenching or smaller pion masses
- Slope in q^2 is consistent with previous result, but with higher preliminary statistical error
- Chromo-magnetic operator is work-in-progress

Thank you!

-  D. Becirevic, V. Lubicz, G. Martinelli and F. Mescia [SPQcdR Collaboration], Phys. Lett. B **501**, 98 (2001) [arXiv:hep-ph/0010349].
-  P. Boucaud *et al.* [ETM Collaboration], Phys. Lett. B **650**, 304 (2007) [arXiv:hep-lat/0701012].
-  R. Baron *et al.* [ETM Collaboration], arXiv:0911.5061 [hep-lat].
-  M. Constantinou *et al.*, arXiv:1004.1115 [hep-lat].