

## Conclusion

# Mathematical Structure of QCD: Wilson Coefficients and Anomalous Dimensions

Johannes Blümlein (1887-1925) was a German chemist who made significant contributions to the field of organic chemistry, particularly in the area of heterocyclic compounds. He is best known for his work on the synthesis of purine nucleotides, which laid the foundation for the development of antibiotics like streptomycin.

view of the most follows in how this model can be used to predict the evolution of the system. The first step is to consider the effect of the initial conditions on the evolution of the system. This can be done by solving the equations of motion for different initial conditions and comparing the results. The second step is to compare the results obtained from the numerical simulations with the analytical predictions. This can be done by plotting the results and comparing them. The third step is to analyze the results and draw conclusions about the behavior of the system. This can be done by examining the results and identifying any trends or patterns. The fourth step is to use the results to make predictions about the future behavior of the system. This can be done by extrapolating the results and using them to predict the future behavior of the system.

# Multiple Harmonic Sums to Level 6

### 3. Multiple Harmonic Sums to Level 6

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#### **4. A Quadratic Law?**

## **5. The 16th New Single Member**

## 5. The 16th Non-Singlet Moment

## 6. Conclusions

It is important that each community member be fully informed about the proposed changes as early as possible to facilitate smooth implementation.

# 1. Introduction

Consider hard scattering processes in massless field theories:

QCD, QED,  $m_i \rightarrow 0$

Factorization Theorem Leading Twist:

The cross section  $\sigma$  factorizes as

$$\sigma = \sum_k \sigma_{k,W} \otimes f_k$$

$\sigma_W$  perturbative Wilson Coefficient

$f$  non-perturbative Parton Density

$\otimes$  Mellin convolution

$$[A \otimes B](x) = \int_0^1 dx_1 \int_0^1 dx_2 \delta(x - x_1 x_2) A(x_1) B(x_2)$$

$$\mathbf{M}[A \otimes B](N) = \mathbf{M}[A](N) \cdot \mathbf{M}[B](N)$$

with the Mellin transform :

$$\mathbf{M}[f(x)](N) = \int_0^1 dx x^{N-1} f(x), \quad \text{Re}[N] > c$$

**Observation :**

Feynman Amplitudes seem to obey the **Mellin Symmetry**

i.e. to significantly simplify in **Mellin Space**

## 2. $x$ Space Results

Usual Starting Point of Higher Order Calculations :

⇒ Nielsen type Integrals and their Generalization

$$S_{n,p,q}(x) = \frac{(-1)^{n+p+q-1}}{(n-1)!p!q!} \int_0^1 \frac{dz}{z} \ln^{(n-1)}(z) \ln^p(1-zx) \ln^q(1+zx)$$

Special Cases:

$$\text{Li}_n(x) = S_{n-1,1}(x) \quad w = n$$

$$\frac{d\text{Li}_2(\pm x)}{d \ln(x)} = -\ln(1 \mp x) \quad w = 1$$

$$\text{Li}_0(x) = \frac{x}{1-x} \quad w = 0$$

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## van Neerven, Zijlstra 1992

$$\begin{aligned}
c_{2,-}^{(2)}(x) = & C_F (C_F - C_A/2) \times \\
& \left\{ \frac{1+x^2}{1-x} \left[ [4 \ln^2(x) - 16 \ln(x) \ln(1+x) - 16 \text{Li}_2(-x) - 8 \zeta_2] \ln(1-x) \right. \right. \\
& + [-2 \ln^2(x) + 20 \ln(x) \ln(1+x) - 8 \ln^2(1+x) + 8 \text{Li}_2(1-x) + 16 \text{Li}_2(-x) - 8] \ln(x) \\
& - 16 \ln(1+x) \text{Li}_2(-x) - 8 \zeta_2 \ln(1+x) - 16 \left[ \text{Li}_3\left(-\frac{1-x}{1+x}\right) - \text{Li}_3\left(\frac{1-x}{1+x}\right) \right] \\
& \left. \left. - 16 \text{Li}_2(1-x) + 8 S_{1,2}(1-x) + 8 \text{Li}_3(-x) - 16 S_{1,2}(-x) + 8 \zeta_3 \right] \right. \\
& +(4+20x) \left[ \ln^2(x) \ln(1+x) - 2 \ln(x) \ln^2(1+x) - 2 \zeta_2 \ln(1+x) - 4 \ln(1+x) \text{Li}_2(-x) \right. \\
& \left. + 2 \text{Li}_3(-x) - 4 S_{1,2}(-x) + 2 \zeta_3 \right] + \left( 32 + 32x + 48x^2 - \frac{72}{5}x^3 + \frac{8}{5x^2} \right) \\
& \times [\text{Li}_2(-x) + \ln(x) \ln(1+x)] + 8(1+x) [\text{Li}_S(1-x) + \ln(x) \ln(1-x)] + 16(1-x) \ln(1-x) \\
& + \left( -4 - 16x - 24x^2 + \frac{36}{5}x^3 \right) \ln^2(x) + \frac{1}{5} \left( -26 - 106x + 72x^2 - \frac{8}{x} \right) \ln(x) \\
& \left. + \left( -4 + 20x + 48x^2 - \frac{72}{5}x^3 \right) \zeta_2 + \frac{1}{5} \left( -162 + 82x + 72x^2 + \frac{8}{x} \right) \right\}
\end{aligned}$$

.... several other pages for  $c_2^{(+)}(x), c_2^G(x), c_L^{(q,G)}(x)$

⇒ 77 Functions @ 2 Loops

⇒ partly rather complicated arguments

⇒ relations are not directly visible ...

The 77 functions do roughly correspond in number to the number of all possible harmonic sums up to weight w=4: 80.

## x Space Results

No.	$f(z)$	$M[f](N) = \int_0^1 dz z^{N-1} f(z)$
1	$\delta(1 - z)$	1
2	$z^r$	$\frac{1}{N + r}$
3	$\left(\frac{1}{1-z}\right)_+$	$-S_1(N-1)$
4	$\frac{1}{1+z}$	$(-1)^{N-1} [\log(2) - S_1(N-1)] + \frac{1 + (-1)^{N-1}}{2} S_1\left(\frac{N-1}{2}\right) - \frac{1 - (-1)^{N-1}}{2} S_1\left(\frac{N-2}{2}\right)$
5	$z^r \log^n(z)$	$\frac{(-1)^n}{(N+r)^{n+1}} \Gamma(n+1)$
6	$z^r \log(1-z)$	$-\frac{S_1(N+r)}{N+r}$
7	$z^r \log^2(1-z)$	$\frac{S_1^2(N+r) + S_2(N+r)}{N+r}$
8	$z^r \log^3(1-z)$	$-\frac{S_1^3(N+r) + 3S_1(N+r)S_2(N+r) + 2S_3(N+r)}{N+r}$
9	$\left[\frac{\log(1-z)}{1-z}\right]_+$	$\frac{1}{2} S_1^2(N-1) + \frac{1}{2} S_2(N-1)$
10	$\left[\frac{\log^2(1-z)}{1-z}\right]_+$	$-\left[\frac{1}{3} S_1^3(N-1) + S_1(N-1)S_2(N-1) + \frac{2}{3} S_3(N-1)\right]$
11	$\left[\frac{\log^3(1-z)}{1-z}\right]_+$	$\frac{1}{4} S_1^4(N-1) + \frac{3}{2} S_1^2(N-1)S_2(N-1) + \frac{3}{4} S_2^2(N-1) + 2S_1(N-1)S_3(N-1) + \frac{3}{2} S_4(N-1)$
12	$\frac{\log^n(z)}{1-z}$	$(-1)^{n+1} \Gamma(n+1) [S_{n+1}(N-1) - \zeta(n+1)]$

Only single sums.

No.	$f(z)$	$M[f](N)$
64	$\frac{\text{Li}_3(-z)}{1+z}$	$(-1)^{N-1} \left\{ S_{3,-1}(N-1) + [S_3(N-1) - S_{-3}(N-1)] \log 2 \right. \\ \left. + \frac{1}{2} \zeta(2) S_{-2}(N-1) - \frac{3}{4} \zeta(3) S_{-1}(N-1) \right. \\ \left. + \frac{1}{8} \zeta^2(2) - \frac{3}{4} \zeta(3) \log 2 \right\}$
65	$\text{Li}_3(1-z)$	$\frac{1}{N} [S_1(N) S_2(N) - \zeta(2) S_1(N) + S_3(N) \\ - S_{2,1}(N) + \zeta(3)]$
66	$\frac{\text{Li}_3(1-z)}{1-z}$	$-S_{1,1,2}(N-1) + \frac{1}{2} \zeta(2) S_1^2(N-1) + \frac{1}{2} \zeta(2) S_2(N-1) \\ - \zeta(3) S_1(N-1) + \frac{2}{5} \zeta^2(2)$
67	$\frac{\text{Li}_3(1-z)}{1+z}$	$(-1)^{N-1} \left[ S_{-1,1,2}(N-1) - \zeta(2) S_{-1,1}(N-1) \right. \\ \left. + \zeta(3) S_{-1}(N-1) + \text{Li}_4\left(\frac{1}{2}\right) - \frac{9}{20} \zeta^2(2) \right. \\ \left. + \frac{7}{8} \zeta(3) \log 2 + \frac{1}{2} \zeta(2) \log^2 2 + \frac{1}{24} \log^4 2 \right]$
68	$\text{Li}_3\left(\frac{1-z}{1+z}\right) \\ - \text{Li}_3\left(-\frac{1-z}{1+z}\right)$	$\frac{(-1)^N}{N} \left[ -S_{-1,2}(N) - S_{-2,1}(N) + S_1(N) S_{-2}(N) \right. \\ \left. + S_{-3}(N) \right. \\ \left. + \zeta(2) S_{-1}(N) + \frac{1}{2} \zeta(2) S_1(N) - \frac{7}{8} \zeta(3) + \frac{3}{2} \zeta(2) \log 2 \right] \\ \left. + \frac{1}{N} \left[ -S_{-1,-2}(N) - S_{2,1}(N) + S_1(N) S_2(N) + S_3(N) \right. \right. \\ \left. \left. - \frac{1}{2} \zeta(2) S_{-1}(N) - \zeta(2) S_1(N) + \frac{21}{8} \zeta(3) - \frac{3}{2} \zeta(2) \log 2 \right] \right]$
69	$\frac{1}{1+z} \left[ \text{Li}_3\left(\frac{1-z}{1+z}\right) \\ - \text{Li}_3\left(-\frac{1-z}{1+z}\right) \right]$	$(-1)^{N-1} \left\{ \text{S}_{1,1,-2}(N-1) - \text{S}_{1,-1,2}(N-1) \right. \\ \left. + \text{S}_{-1,1,2}(N-1) - \text{S}_{-1,-1,-2}(N-1) \right. \\ \left. + 2\zeta(2) S_{1,-1}(N-1) + \frac{1}{4} \zeta(2) S_1^2(N-1) - \frac{1}{4} \zeta(2) S_{-1}^2(N-1) \right. \\ \left. - \zeta(2) S_1(N-1) S_{-1}(N-1) - \zeta(2) S_{-2}(N-1) \right. \\ \left. - \left[ \frac{7}{8} \zeta(3) - \frac{3}{2} \zeta(2) \log 2 \right] S_1(N-1) \right. \\ \left. + \left[ \frac{21}{8} \zeta(3) - \frac{3}{2} \zeta(2) \log 2 \right] S_{-1}(N-1) \right. \\ \left. - 2\text{Li}_4\left(\frac{1}{2}\right) + \frac{19}{40} \zeta^2(2) + \frac{1}{2} \zeta(2) \log^2 2 - \frac{1}{12} \log^4 2 \right\}$

2 loop coefficient functions  $\Rightarrow$  Nested Harmonic Sums of  
Weight  $w = 4$

## x Space Results

$$\begin{aligned}
 S_{-1,-1,-2}(N) = & \\
 (-1)^{N+1} \mathbf{M} \left\{ \frac{1}{1+x} [F_1(x) + \log(1-x)\text{Li}_2(-x)] \right\} (N) & \\
 + (-1)^{N+1} \mathbf{M} \left\{ \frac{1}{1+x} \left[ \frac{1}{2} S_{1,2}(x^2) - S_{1,2}(x) - S_{1,2}(-x) \right] \right\} (N) & \\
 + \frac{1}{2} \zeta(2) [S_{-1,1}(N) - S_{-1,-1}(N)] & \\
 + \left[ \frac{9}{8} \zeta(3) - \frac{3}{2} \zeta(2) \log(2) - \frac{1}{6} \log^3(2) \right] S_{-1}(N) & \\
 - \frac{1}{10} \zeta(2)^2 + \frac{17}{8} \zeta(3) \log(2) - \frac{7}{4} \zeta(2) \log^2(2) - \frac{1}{6} \log^4(2) &
 \end{aligned}$$

with

$$\begin{aligned}
 F_1(x) = & S_{1,2} \left( \frac{1-x}{2} \right) + S_{1,2}(1-x) - S_{1,2} \left( \frac{1-x}{1+x} \right) \\
 & + S_{1,2} \left( \frac{1}{1+x} \right) - \ln(2) \left( \frac{1-x}{2} \right) \\
 & + \frac{1}{2} \ln^2(2) \ln \left( \frac{1+x}{2} \right) - \ln(2) \text{Li}_2 \left( \frac{1-x}{1+x} \right)
 \end{aligned}$$

$F_1(x)$ , although of complicated structure, it reduces completely via algebraic relations

⇒ Mellin polynomial of simpler objects

These objects can be very complicated integrals. J.B., van Neerven, Ravindran, Kawamura 2000, 2003

### 3. Multiple Harmonic Sums to Level 6

The simplest example :

$$P_{qq}(x) = \left( \frac{1+x^2}{1-x} \right)_+ = \frac{2}{(1-x)_+} + \dots$$

$$\int_0^1 dx \frac{x^{N-1}}{(1-x)_+} = - \sum_{k=0}^{N-2} \int_0^1 dx x^k = - \sum_{k=1}^{N-1} \frac{1}{k} = -S_1(N-1)$$

Alternating sums :

$$S_{-1}(N-1) = (-1)^{N-1} \mathbf{M} \left[ \frac{1}{1+x} \right] (N) - \ln(2) = \int_0^1 dx \frac{x^{N-1}}{(1-x)_+} = \sum_{k=1}^{N-1} \frac{(-1)^k}{k}$$

(Finite for  $N \rightarrow \infty$ .)

General case :

$$S_{a_1, \dots, a_l}(N) = \sum_{k_1=1}^N \frac{(\text{sign}(a_1))^{k_1}}{k_1^{\|a_1\|}} \sum_{k_2=1}^N \frac{(\text{sign}(a_2))^{k_2}}{k_2^{\|a_2\|}} \dots$$

Vermaseren, 1997

All Mellin transforms occurring in massless Field Theories for 1-Parameter Quantities can be represented by Harmonic Sums (at least to 3-loop order).

## Algebraic Relations

First relation:

L. Euler, 1775

$$S_{m,n} + S_{n,m} = S_m \cdot S_n + S_{m+n}, \quad m, n > 0$$

Generalized to alternating sums by

$$\begin{aligned} S_{m,n} + S_{n,m} &= S_m \cdot S_n + S_{m \wedge n}, \\ m \wedge n &= [|m| + |n|] \text{sign}(m)\text{sign}(n) \end{aligned}$$

Ternary relations: Sita Ramachandra Rao, 1984,

4-ary relation: J.B., Kurth, 1998.

These & other relations hold widely independent  
of their **Value** and **Type**.

Determined by : • Index Structure  
• Multiplication Relation

The Formalism applies as well to the Harmonic Polylogarithms.

Remiddi, Vermaseren, 1999.

Application to QED: T. Riemann et al., 2004

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## Linear Representations of Mellin Transform by Harmonic Sums:

$$\mathbf{M}[F_w(x)](N) = S_{k_1, \dots, k_m}^w(N) + P\left(S_{k_1, \dots, k_r}^{\tau'}, \sigma_{k_1, \dots, k_p}^{\tau''}\right)$$

$$w = \sum_{i=1}^m \|k_i\| \quad \text{Weight}$$

$$\tau', \tau'' < w \quad P \text{ is a polynomial.}$$

w	#	$\Sigma$	
1	2	2	
2	6	8	
3	18	26	2 Loop anom. Dimensions
4	54	80	2 Loop Wilson Coefficients
5	162	242	3 Loop anom. Dimensions
6	486	728	3 Loop Wilson Coefficients
$2 \cdot 3^{w-1}$		$3^w - 1$	

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## Shuffle Products

Depth 2:

$$S_{a_1}(N) \sqcup\sqcup S_{a_2}(N) = S_{a_1, a_2}(N) + S_{a_2, a_1}(N)$$

Depth 3:

$$S_{a_1}(N) \sqcup\sqcup S_{a_2, a_3}(N) = S_{a_1, a_2, a_3}(N) + S_{a_2, a_1, a_3}(N) + S_{a_2, a_3, a_1}(N)$$

Depth 4:

$$\begin{aligned} S_{a_1}(N) \sqcup\sqcup S_{a_2, a_3, a_4}(N) &= S_{a_1, a_2, a_3, a_4}(N) + S_{a_2, a_1, a_3, a_4}(N) + S_{a_2, a_3, a_1, a_4}(N) \\ &+ S_{a_2, a_3, a_4, a_1}(N) \\ S_{a_1, a_2}(N) \sqcup\sqcup S_{a_3, a_4}(N) &= S_{a_1, a_2, a_3, a_4}(N) + S_{a_1, a_3, a_2, a_4}(N) + S_{a_1, a_3, a_4, a_2}(N) \\ &+ S_{a_3, a_4, a_1, a_2}(N) + S_{a_3, a_1, a_4, a_2}(N) + S_{a_3, a_1, a_2, a_4}(N) \end{aligned}$$

Depth 5:

$$\begin{aligned} S_{a_1}(N) \sqcup\sqcup S_{a_2, a_3, a_4, a_5}(N) &= S_{a_1, a_2, a_3, a_4, a_5}(N) + S_{a_2, a_1, a_3, a_4, a_5}(N) \\ &+ S_{a_2, a_3, a_1, a_4, a_5}(N) + S_{a_2, a_3, a_4, a_1, a_5}(N) \\ &+ S_{a_2, a_3, a_4, a_5, a_1}(N) \\ S_{a_1, a_2}(N) \sqcup\sqcup S_{a_3, a_4, a_5}(N) &= S_{a_1, a_2, a_3, a_4, a_5}(N) + S_{a_1, a_3, a_2, a_4, a_5}(N) \\ &+ S_{a_1, a_3, a_4, a_2, a_5}(N) + S_{a_1, a_3, a_4, a_5, a_2}(N) \\ &+ S_{a_3, a_1, a_2, a_4, a_5}(N) + S_{a_3, a_1, a_4, a_2, a_5}(N) \\ &+ S_{a_3, a_1, a_4, a_5, a_2}(N) + S_{a_3, a_4, a_5, a_1, a_2}(N) \\ &+ S_{a_3, a_4, a_1, a_5, a_2}(N) + S_{a_3, a_4, a_1, a_2, a_5}(N) \end{aligned}$$

Depth 6: .....

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## Algebraic Equations

Depth 2:

$$S_{a_1}(N) \sqcup\!\!\! \sqcup S_{a_2}(N) - S_{a_1}(N)S_{a_2}(N) - S_{a_1 \wedge a_2}(N) = 0$$

Depth 3:

$$S_{a_1}(N) \sqcup\!\!\! \sqcup S_{a_2, a_3}(N) - S_{a_1}(N)S_{a_2, a_3}(N) - S_{a_1 \wedge a_2, a_3}(N) - S_{a_2, a_1 \wedge a_3}(N) = 0$$

Depth 4:

$$\begin{aligned} S_{a_1}(N) \sqcup\!\!\! \sqcup S_{a_2, a_3, a_4}(N) & - S_{a_1}(N)S_{a_2, a_3, a_4}(N) - S_{a_1 \wedge a_2, a_3, a_4}(N) \\ & - S_{a_2, a_1 \wedge a_3, a_4}(N) - S_{a_2, a_3, a_1 \wedge a_4}(N) = 0 \\ S_{a_1, a_2}(N) \sqcup\!\!\! \sqcup S_{a_3, a_4}(N) & - S_{a_1, a_2}(N)S_{a_3, a_4}(N) - S_{a_1, a_2 \wedge a_3, a_4}(N) \\ & - S_{a_1, a_3, a_2 \wedge a_4}(N) - S_{a_3, a_1 \wedge a_4, a_2}(N) \\ & - S_{a_3, a_1, a_2 \wedge a_4}(N) - S_{a_1 \wedge a_3, a_2, a_4}(N) \\ & - S_{a_1 \wedge a_3, a_4, a_2} + S_{a_1 \wedge a_3, a_2 \wedge a_4} = 0 \end{aligned}$$

Depth 5: .....

# Basic Sums = # Permutations - # Independent Equations

## Some Solution for $d = 6$

$$\begin{aligned}
S_{a,a,a,a,b,b} = & \\
& - \frac{1}{4} S_a S_b, a, a, a, b + \frac{3}{4} S_a \wedge b, a, a, a, b - \frac{1}{4} S_{b,a,a,a,a \wedge b} + \frac{1}{12} S_a S_{a,a,b,b,a} + S_{a,a,a,a,b \wedge b} \\
& - \frac{1}{12} S_{a \wedge a,b,b,a,a} - \frac{1}{12} S_{a,b,b,a \wedge a,a} - \frac{1}{12} S_{a,b,b,a,a \wedge a} - \frac{1}{4} S_{b,a \wedge a,a,a,b} - \frac{1}{4} S_{b,a,a \wedge a,a,b} \\
& - \frac{1}{4} S_{b,a,a,a \wedge a,b} - \frac{1}{4} S_{a,a \wedge a,a,b,b} - \frac{1}{4} S_{a,a,a \wedge a,b,b} - \frac{1}{4} S_{a,b,a \wedge a,a,b} - \frac{1}{4} S_{a,b,a,a \wedge a,b} \\
& + \frac{1}{12} S_{a \wedge a,b,a,b,a} + \frac{1}{12} S_{a,b,a \wedge a,b,a} - \frac{1}{4} S_{a,a,a,b,a \wedge b} - \frac{1}{4} S_{a,a,b,a,a \wedge b} + \frac{1}{12} S_{a,a,a \wedge b,b,a} \\
& + \frac{3}{4} S_{a,a,a \wedge b,a,b} - S_{b,b,a,a,a,a} + \frac{1}{4} S_{b,b,a,a \wedge a,a} - \frac{1}{4} S_{a \wedge a,a,a,b,b} + \frac{1}{12} S_{a,b,a,b,a \wedge a} \\
& + \frac{1}{12} S_{b,a \wedge a,a,b,a} + \frac{1}{12} S_{b,a,a \wedge a,b,a} + \frac{1}{12} S_{b,a,a,b,a \wedge a} - \frac{1}{12} S_{b,a \wedge a,b,a,a} - \frac{1}{12} S_{b,a,b,a \wedge a} \\
& + \frac{1}{4} S_{b,b,a \wedge a,a,a} + \frac{1}{4} S_{b,b,a,a,a \wedge a} - \frac{1}{4} S_{a \wedge a,a,b,a,b} - \frac{1}{4} S_{a,a \wedge a,b,a,b} - \frac{1}{4} S_{a,a,b,a \wedge a,b} \\
& + \frac{1}{12} S_{a \wedge a,a,b,b,a} + \frac{1}{12} S_{a,a \wedge a,b,b,a} + \frac{1}{12} S_{a,a,b,b,a \wedge a} - \frac{1}{4} S_{a \wedge a,b,a,a,b} - \frac{1}{12} S_{b,a,b,a,a \wedge a} \\
& + \frac{1}{12} S_{a \wedge b,a,a,b,a} + \frac{1}{12} S_{b,a,a,a \wedge b,a} - \frac{1}{12} S_{a \wedge b,a,b,a,a} + \frac{1}{4} S_{a \wedge b,b,a,a,a} + \frac{1}{4} S_{b,a \wedge b,a,a,a} \\
& - \frac{1}{12} S_{b,a,a \wedge b,a,a} + \frac{3}{4} S_{a,a,a,a \wedge b,b} + \frac{1}{12} S_{a,a,b,a \wedge b,a} + \frac{3}{4} S_{a,a \wedge b,a,a,b} - \frac{1}{4} S_{a,b,a,a,a \wedge b} \\
& + \frac{1}{12} S_{a,a \wedge b,a,b,a} + \frac{1}{12} S_{a,b,a,a \wedge b,a} - \frac{1}{12} S_{a,a \wedge b,b,a,a} - \frac{1}{12} S_{a,b,a,a,a,b} - \frac{1}{4} S_a S_{a,a,a,a,b}, \\
& - \frac{1}{12} S_a S_{a,b,b,a,a} + \frac{1}{12} S_a S_{a,b,a,b,a} - \frac{1}{12} S_a S_{b,a,b,a,a} + \frac{1}{12} S_a S_{b,a,a,b,a} - \frac{1}{4} S_a S_{a,b,a,a,a}, \\
& + S_b S_{a,a,a,a,b} + \frac{1}{4} S_a S_{b,b,a,a,a} - \frac{1}{4} S_a S_{a,a,b,a,b}
\end{aligned}$$

## Depth $d = 3$

Index Set	Number	Dep. Sums of Depth 3	min. Weight	Fraction of fundl. Sums
$\{a, a, a\}$	1	1	3	0
$\{a, a, b\}$	3	2	3	$1/3$
$\{a, b, c\}$	6	4	4	$1/3$

## Depth $d = 4$

Index Set	Number	Dep. Sums of Depth 4	min. Weight	Fraction of fundl. Sums
$\{a, a, a, a\}$	1	1	4	0
$\{a, a, a, b\}$	4	3	4	$1/4$
$\{a, a, b, b\}$	6	5	4	$1/6$
$\{a, a, b, c\}$	12	9	5	$1/4$
$\{a, b, c, d\}$	24	18	6	$1/4$

## Depth $d = 6$

Index Set	Number	Dep. Sums of Depth 6	min. Weight	Fraction of fundl. Sums
$\{a, a, a, a, a, a\}$	1	1	6	0
$\{a, a, a, a, a, b\}$	6	5	6	$1/6$
$\{a, a, a, a, b, b\}$	15	13	6	$2/15$
$\{a, a, a, b, b, b\}$	20	17	6	$3/20$
$\{a, a, a, a, b, c\}$	30	25	7	$1/6$
$\{a, a, a, b, b, c\}$	60	50	7	$1/6$
$\{a, a, b, b, c, c\}$	90	76	8	$7/45$
$\{a, a, a, b, c, d\}$	120	100	8	$1/6$
$\{a, a, b, b, c, d\}$	180	150	8	$1/6$
$\{a, a, b, c, d, e\}$	360	300	10	$1/6$
$\{a, b, c, d, e, f\}$	720	600	12	$1/6$

# Theory of Words

Can we count the Basis in simpler way ?  $\Rightarrow$  YES.

Free Algebras and Elements of the Theory of Codes  
 $\Rightarrow$  Particle Physics

Only the multiplication relation  
and the Index structure matters

$\mathfrak{A} = \{a, b, c, d, \dots\}$       Alphabet

$a < b < c < d < \dots$       ordered

$\mathfrak{A}^*(\mathfrak{A})$  Set of all words **W**

$W = a_1 \cdot a_2 \cdot a_3 \dots a_{532} \equiv$  concatenation product (nc)

$W = p \cdot x \cdot s$     p = prefix; s = suffix

Definition:

A Lyndon word is smaller than any of its suffixes.

Theorem: [Radford, 1979]

The shuffle algebra  $K\langle\mathfrak{A}\rangle$  is freely generated by the Lyndon words.  
I.e. the number of Lyndon words yields the number of basic elements.

Examples :

$\{a, a, \dots, a, b\} = aaa \dots ab$       1 Lyndon word for these sets

$n \rightarrow a's : n_{basic}/n_{all} = 1/n$        $n \equiv$  depth of the sums

$\{a, a, a, b, b, b\}$     *aaabbb, aababb, aabbab*    3 Lyndon words

$n_{basic}/n_{all} = 3/20 < 1/6$ . Symmetries lead to a smaller fraction.

## Is there a general Counting Relation ?

E. Witt, 1937

$$l_n(n_1, \dots, n_q) = \frac{1}{n} \sum_{d \mid n_i} \mu(d) \frac{(n/d)!}{(n_1/d)! \dots (n_q/d)!}, \quad \sum_i n_i = n$$

$\mu(k)$  Möbius function

2nd Witt formula.

The Length of the Basis is a function mainly of the Depth.

$$l_6(\{a, a, a, b, b, b\}) = \frac{1}{6} \left[ \mu(1) \frac{6!}{3!3!} + \mu(3) \frac{2!}{1!1!} \right] = 3$$

$$n_6(\{a, a, a, b, b, b\}) = \frac{6!}{2!3!} = 20$$

Weight	# Sums	Cum. # Sums	# Basic Sums	Cum. # Basic Sums	Cum. Fraction
1	2	2	0	0	0.0
2	6	8	1	1	0.1250
3	18	26	6	7	0.2692
4	54	80	16	23	0.2875
5	162	242	46	69	0.2851
6	486	728	114	183	0.2513

↑ 2nd Witt formula

## Structural Relations

Seek for further Reduction:

Relations using the Value of the Objects [DESY 04-064]

Use Relations like:

$$\frac{1}{2} \frac{\text{Li}_2(x^2)}{1-x^2} = \frac{\text{Li}_2(x)}{1-x} + \frac{\text{Li}_2(x)}{1+x} + \frac{\text{Li}_2(-x)}{1-x} + \frac{\text{Li}_2(-x)}{1+x}$$

and similar ones.

Apply Symmetries among Mellin-Transforms of Nielsen Integrals.

Since all harmonic sums are meromorphic functions for  $N \in \mathbb{C}$  since they may be represented by Factorial Series Derivatives are not essentially new functions.

$$\mathbf{M}[\ln^k(x)f(x)](N) = \frac{\partial^k}{\partial N^k} \mathbf{M}[f(x)](N)$$

Further Reduction due to the Structure of Feynman Amplitudes

The Lord is mercy, after all!!

## Analytic Continuation

The Harmonic Sums and Mellin Transforms have to be represented such, that the outer summation index can be analytically continued to  $N \in \mathbb{C}$  [J.B., 2000]

- Use precise, adaptive Representations in analytic Form
- Refer to the Representation through Factorial Series etc.
- The **Residue Theorem** is used to get back to  $x$  space

## 4. A Quadratic Law ?

The anomalous dimensions and Wilson coefficients for  $m_i = 0$  can be expressed in terms of multiple harmonic sums to 3-loop order.

What are the irreducible functions behind this representation ?

We will not count Euler's  $\Gamma$ -function neither all derivations of the functions occurring.

### The final set of functions:

#### Trivial functions:

$$S_{\pm k}(N) \longrightarrow \psi^{(k-1)}(N+1)$$

For  $w = 1, 2$  no non-trivial functions contribute to the anomalous dimensions and Wilson coefficients.

#### Non-trivial functions:

$N = 3$  : Two-Loop anomalous dimensions

$$\mathbf{M} \left[ \frac{\text{Li}_2(x)}{1+x} \right] (N)$$

Yndurain et al., 1980

$N = 4$  : Two-Loop Wilson Coefficients

$$\mathbf{M} \left[ \frac{\text{Li}_2(x)}{1-x} \right] (N), \quad \mathbf{M} \left[ \frac{\text{Li}_3(x)}{1+x} \right] (N), \quad \mathbf{M} \left[ \frac{S_{1,2}(x)}{1 \pm x} \right] (N)$$

J.B., S. Moch, 2003,

also: J.B., V. Ravindran, 2004.

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## $N = 5$ : Three-Loop Anomalous Dimensions

$$\begin{aligned} \mathbf{M} \left[ \frac{\ln(1+x)}{1+x} \right] (N), \quad \mathbf{M} \left[ \frac{\text{Li}_4(x)}{1 \pm x} \right] (N), \quad \mathbf{M} \left[ \frac{S_{1,3}(x)}{1+x} \right] (N), \\ \mathbf{M} \left[ \frac{S_{2,2}(x)}{1 \pm x} \right] (N), \quad \mathbf{M} \left[ \frac{S_{2,2}(-x) - \text{Li}_2^2(-x)/2}{1 \pm x} \right] (N), \\ \mathbf{M} \left[ \frac{\text{Li}_2^2(x)}{1+x} \right] (N) \end{aligned}$$

J.B., S. Moch, 2004.

The number of **Non-trivial Basic Functions** seems to grow as :

$$N_w = \theta(w-2) \cdot [w-2]^2$$

Essentially **14 Functions** seem to rule the single scale processes of massless QCD.

This is a rather small number if compared to the number of possible harmonic sums  $3^w - 1$ .

$$\text{Li}_4(x) = L_{x_0^3 x_1} \quad S_{2,2}(x) = L_{x_0^2 x_1^2}$$

$$S_{1,3}(x) = L_{x_0 x_1^3} \quad \text{Li}_n(x) = L_{x_0^{n-1} x_1}$$

## 5. The 16th of the 3-Loop Non-Singlet Anomalous Dimension of $F_1(x, Q^2)$

Seek for another, blind check of the complete calculation of the NS-anomalous dimension by Moch, Vermaseren, and Vogt, 2004.

The calculation was started far before the complete calculation was completed and is based on the **MINCER** algorithm used before by Larin, Noguiera, van Ritbergen, Vermaseren and Retey, 1994–2000

Moment	CPU time [days]	
	$g_{\mu\nu}$	$P_\mu P_\nu$
2	0.002567	0.002190
4	0.012562	0.020027
6	0.057144	0.059320
8	0.303415	0.332731
10	1.108047	1.219046
.	.	.
16	236.236352	—

J.B., J. Vermaseren, 2004

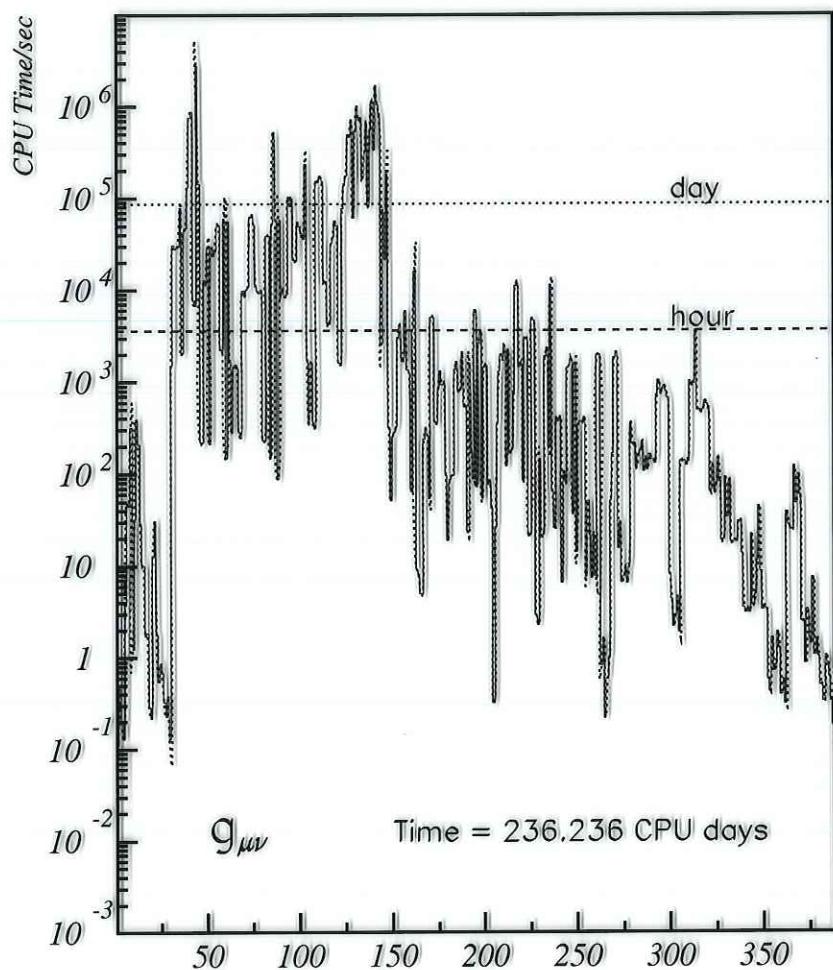
Set-up used: 1 XEON-Dual PC 3 GHz (**Drittmittel**) and 1 XEON-Dual 2.6 GHz PC (borrowed from **AMANDA**) linked to a 4.2 Tbyte RAID system; partly 1 64-bit OPTERON 4Gbyte (dual)

Many Thanks to: P. Wegner, S. Wiesand, U. Gensch & C. Spiering for supporting this project.

Several diagrams stayed in the CPU for 25–40 days each.

→ DESY-Z urgently needs a Parallel PC-Facility ←  
for FORM Formula Manipulation

### Run-time statistics :



Diagram

## Results :

$$\gamma_{16}^{(0)} = \frac{64419601}{6126120} \text{ CF}$$

$$\begin{aligned} \gamma_{16}^{(1)} = & -\frac{1176525373840303}{112588038763200} \text{ CF NF} + \frac{21546159166129889}{484994628518400} \text{ CF CA} \\ & - \frac{3689024452928781382877}{459818557352009856000} \text{ CF}^2 \end{aligned}$$

$$\begin{aligned} \gamma_{16}^{(2)} = & \left( \frac{59290512768143}{1563722760600} \zeta_3 - \frac{58552930270652300886778705063429867}{3451337970612452534317096673280000} \right) \text{ CF}^3 \\ & + \left( -\frac{15018421824060388659436559}{579371382263532418560000} - \frac{64419601}{765765} \zeta_3 \right) \text{ CF CA NF} \\ & + \left( \frac{1670423728083984207878825467}{6488959481351563087872000} + \frac{59290512768143}{3127445521200} \zeta_3 \right) \text{ CF CA}^2 \\ & - \frac{5559466349834573157251}{2069183508084044352000} \text{ CF NF}^2 \\ & + \left( -\frac{1229794646000775781127856064477}{30335885575318557435801600000} - \frac{59290512768143}{1042481840400} \zeta_3 \right) \text{ CF}^2 \text{ CA} \\ & + \left( -\frac{71543599677985155342551355451}{938967886855098206346240000} + \frac{64419601}{765765} \zeta_3 \right) \text{ CF}^2 \text{ NF} \end{aligned}$$

Agreement with : Moch, Vermaseren, Vogt, hep-ph/0403192.

## 7. Conclusions

- Mellin space expressions of anomalous dimensions and Wilson coefficients are of much simpler structure than the  $x$ -space results.
- Index-based algebraic relations of harmonic sums, structural relations of Mellin–transforms of Nielsen–integrals and the specific structure of Feynman amplitudes cause this reduction.
- All algebraic relations are derived in explicit form up to weight  $w = 6$  and apply to harmonic sums, harmonic polylogarithms and all other objects in the corresponding equivalence class. The structural relations are worked out up to  $w = 4$  and the anomalous dimensions for  $w = 5$ .
- The number of multiple harmonic sums of weight  $w$  is  $2 \cdot 3^{w-1}$ . The number of the harmonic sums after algebraic reduction is given by the Witt formula(e) yielding a reduction to  $\approx 1/4$ . Further reductions result from structural relations.
- The number of non-trivial basic functions for  $w \leq 5$  which are needed to express the known anomalous dimensions and the (space– and time–like) Wilson coefficients for  $m_i = 0$  is given by

$$N_w = \theta(w - 2) \cdot [w - 2]^2$$

- An independent calculation of the 16th moment of the non-singlet structure function  $F_1(x, Q^2)$  shows agreement with the complete calculation.