

1

Dijet Photoproduction and Photon Structure at TERA

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and

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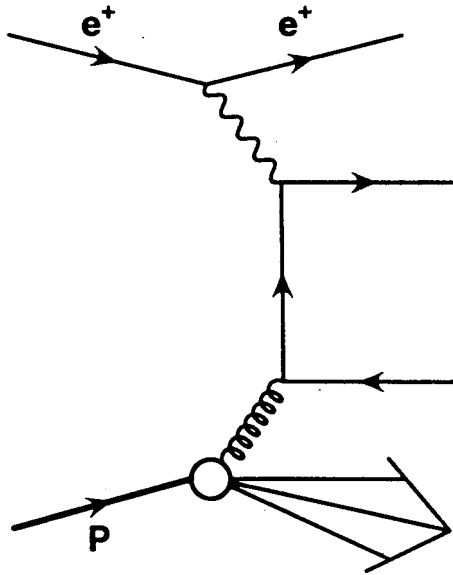
TERA Workshop

9 February 2000

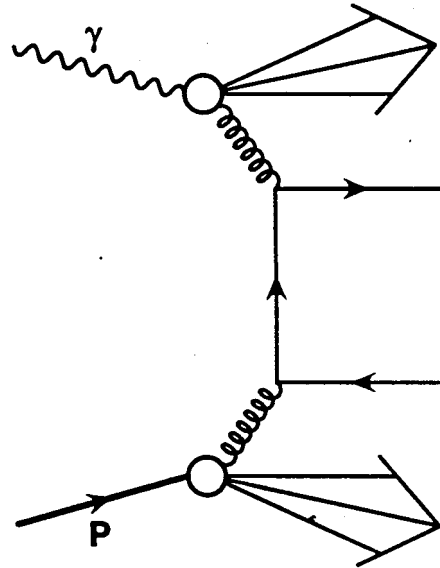
M. Krawczyk

Introduction to Photoproduction

Leading Order (LO) picture:



Direct



Resolved

Resolved cross-section

$$d\sigma_{\gamma p \rightarrow cd} = \sum_{ab} \int_{x_p} \int_{x_\gamma} f_{p \rightarrow b}(x_p, \mu^2) f_{\gamma \rightarrow a}(x_\gamma, \mu^2) \mathcal{M}_{ab \rightarrow cd}$$

$$\mu^2 \Rightarrow \text{scale} \rightarrow (E_T^{\text{jet}})^2$$

$$\mathcal{M}_{ab \rightarrow cd} \Rightarrow \text{perturbatively calculable}$$

$$f_{p \rightarrow b} \Rightarrow \text{experimentally constrained}$$

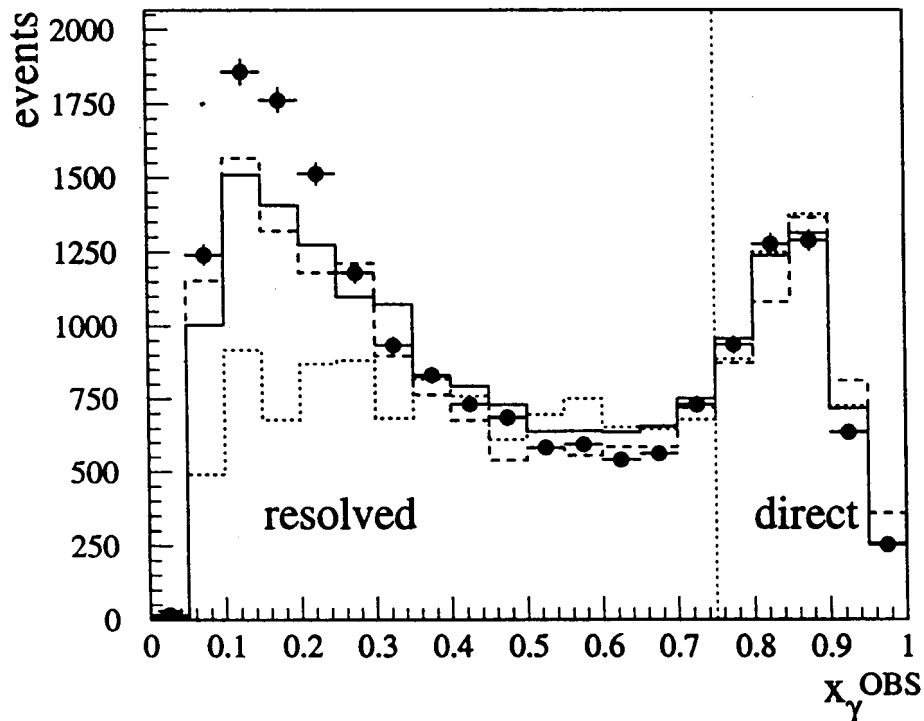
$$f_{\gamma \rightarrow a} \Rightarrow \text{photon structure?}$$

Dijet photoproduction at HERA; lower E_T^{jet}

Here $E_T^{\text{jet}} > 6$ GeV, data compared with MC:

$$x_\gamma^{\text{obs}} = \frac{E_T^{\text{jet1}} e^{-\eta^{\text{jet1}}} + E_T^{\text{jet2}} e^{-\eta^{\text{jet2}}}}{2yE_e}$$

ZEUS 1994



- Large resolved photon cross-section.

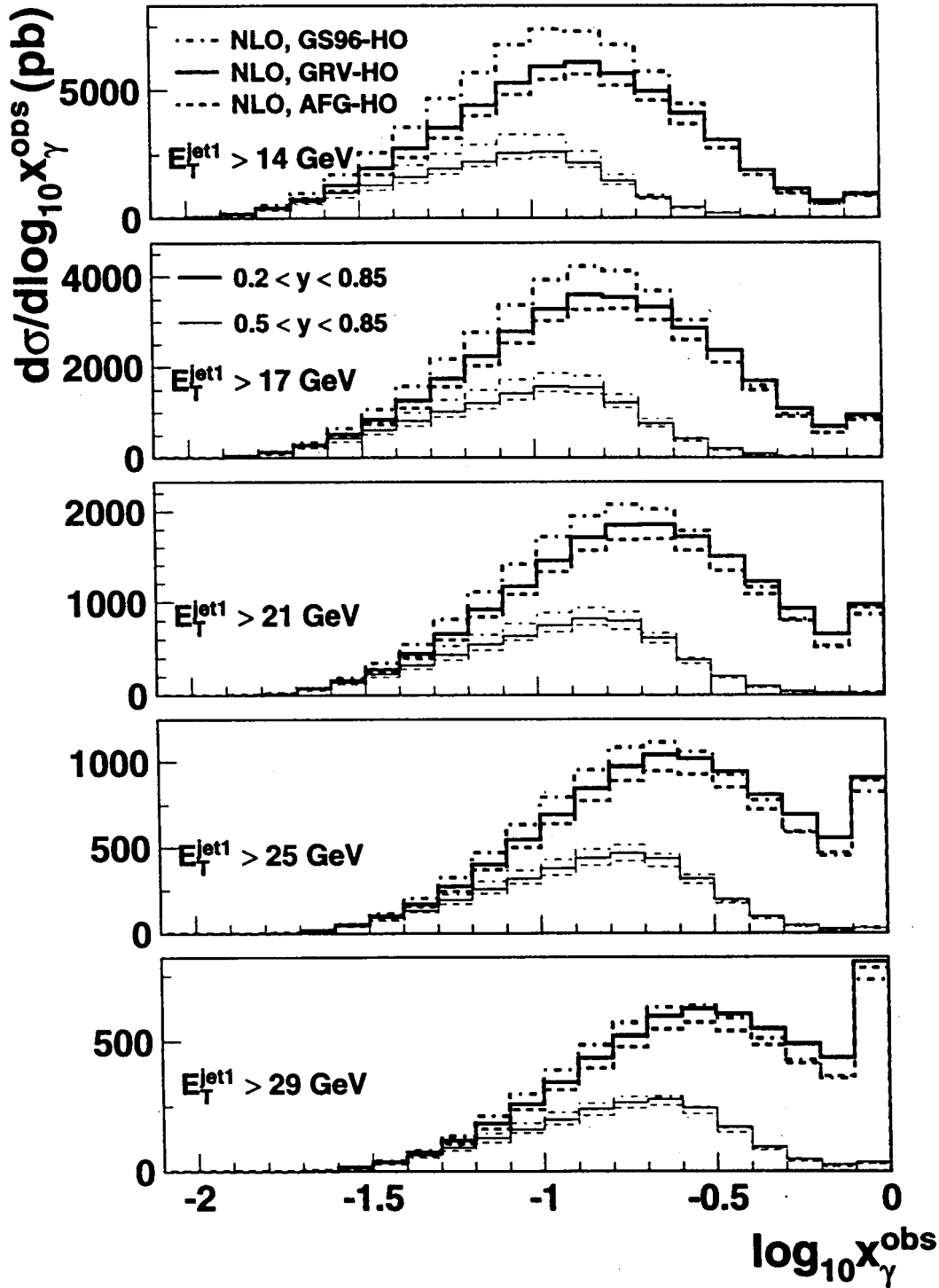
BUT:

- Need additional soft physics model to "describe" data.
- Additional soft physics is obstacle to photon structure sensitivity.
- Answer: goto higher E_T^{jet}

Cross-Sections in x_γ^{obs}

THERA

$$x_\gamma^{\text{obs}} = \frac{E_T^{\text{jet1}} e^{-\eta^{\text{jet1}}} + E_T^{\text{jet2}} e^{-\eta^{\text{jet2}}}}{2yE_e}$$

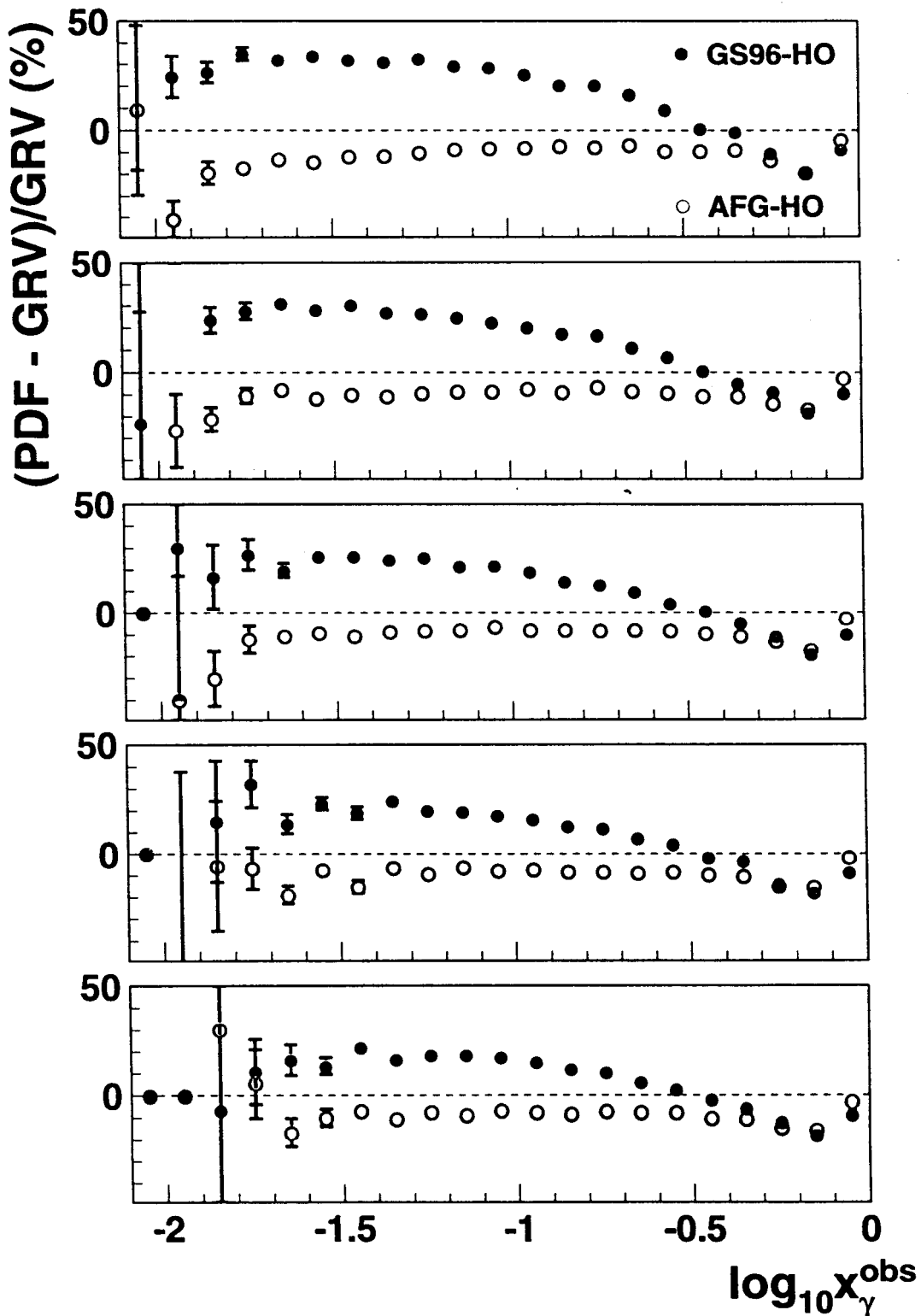


- Large cross-section at low x_γ^{obs} .
- Large variation due to structure function particularly at lower E_T^{jet} .

Differences in PDF

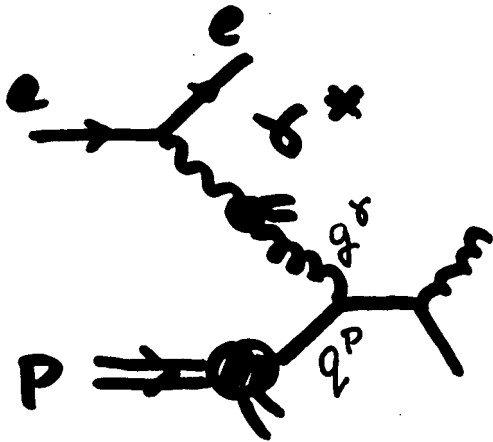
THERA 5

Look at differences with respect to GRV-HO.



- GS96-HO about 30% higher at low x_{γ}^{obs} .
Difference decreasing with increasing x_{γ}^{obs} .
- AFG-HO generally 10% lower than GRV-HO.

PHOTOPRODUCTION at THERA

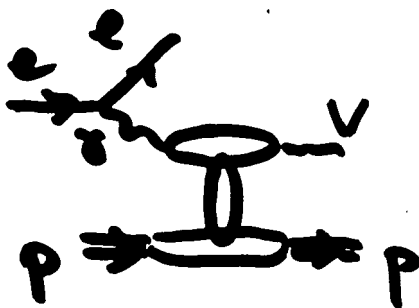


inclusive photoproduction

$\uparrow P_T$

\rightarrow structure of γ and γ^*
(resolved photon processes)

- dijet
- heavy quark
- prompt photon $\hookrightarrow q\gamma^*$



exclusive photoproduction

$\leftarrow \rho, \omega, \dots$

- $\gamma P \rightarrow V P$
- $\gamma P \rightarrow \pi^0, \eta, \eta', \eta_c, \dots, P$ (IG)

J. Butterworth
 M. Wing
S. Söldner-Reubold
 M. Klasen
 B. Surrou
 A. de Roeck
P. Jankowski
 A. Zembruski

P. Landshoff
J. Ginzburg
 M. K.
 F. Cornet
 ...

(Low) Q^2 RANGE AT THERA

7

* PHOTOPRODUCTION AT HERA:

UNTAGGED $Q^2 < 1 \text{ GeV}^2$

$Q^2_{\text{median}} \sim 10^{-3} \text{ GeV}^2$

* AT THERA: (ASSUMING SIMILAR DETECTOR CONFIGURATION)

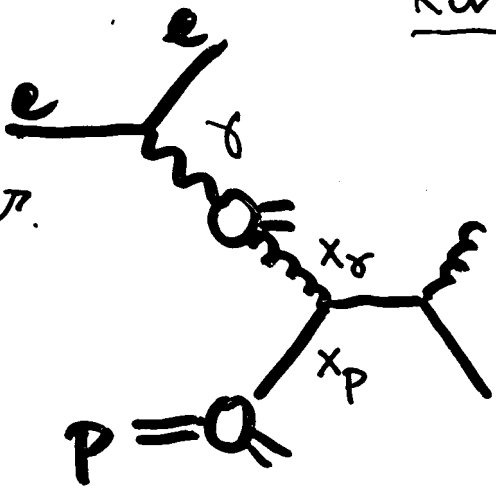
UNTAGGED $Q^2 \lesssim 100 \text{ GeV}^2$

$Q^2_{\text{median}} \sim 10^{-2} \text{ GeV}^2 (?)$

⇒ TWO SCALE PROBLEM!!

* NEED TO INVESTIGATE HAVING A
(LOW Q^2) TAGGER.*

Range of x_γ at THERA



WNA \rightarrow
real γ

$$E_\gamma = \gamma E_e$$

$\uparrow p_T$

$$5 \leq p_T \leq 100 \text{ GeV}$$

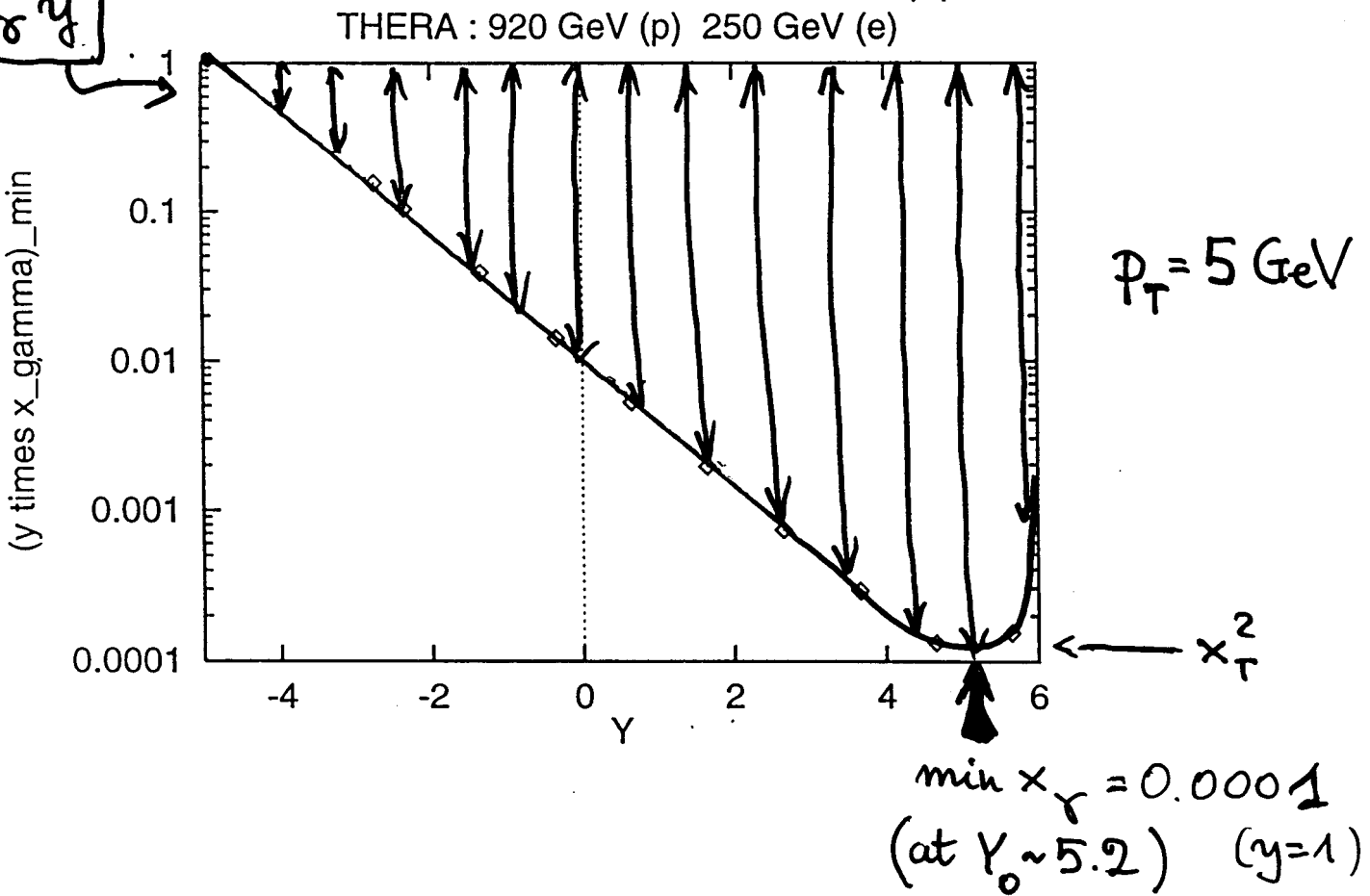
$$-4.6 \leq Y \leq 5.9$$

rapidity range

$$x_p x_\gamma \gamma S_{ep} \geq (2 p_T)^2$$

$$\rightarrow (x_\gamma \gamma)_{\min} = x_T^2 \quad (p_T = 5)$$

range of $x_\gamma \gamma$



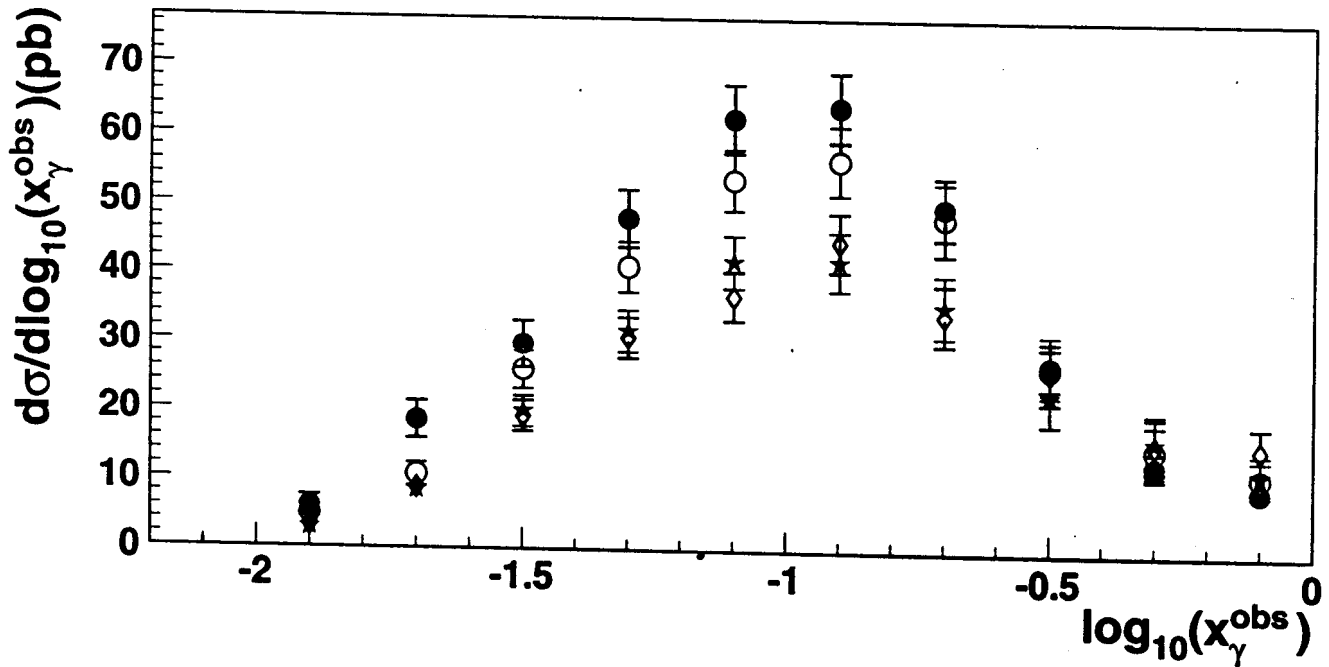
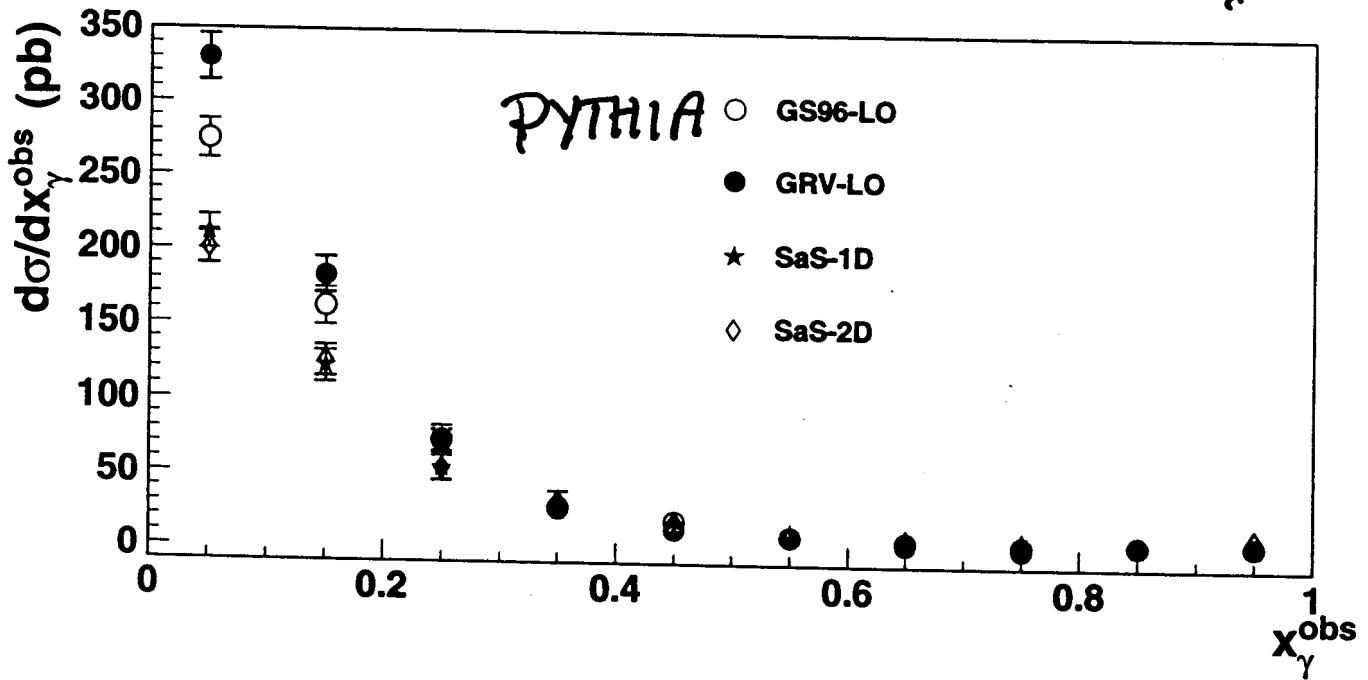
U. Katz $\rightarrow \theta = 0.5^\circ \leftrightarrow$
 $Y = 5.43$
(max exp)

19

Heavy quark production
at THERA

MC $b\bar{b} \rightarrow \mu$ PRODUCTION (M. WING)

PHOTON SENSITIVITY TO CHOICE OF PDF.



* SIMILAR RESULTS FOR HERWIG

* BUT MULTIPARTON INTERACTIONS
COULD PLAY A RÔLE.

LO : c, b, τ

NOWE
P. Janowski

Heavy quark production in TERA, $E_p = 920$ GeV, $E_e = 250$ GeV

77

The calculation was made in LO approximation. Though we used HO parametrization functions:

- photon - GRVHO: M.Glück, E.Reya and A.Vogt, Phys.Rev. D46 (1992) 1973
- proton - MRS(G): A.D.Martin, W.J.Stirling and R.G.Roberts, Phys.Lett. B354 (1995) 155-162

as it was done in following work:

B.A.Kniehl, M.Krämer, G.Kramer and M.Spira, Phys. Lett. B356 (1995) 539-545; hep-ph/9505410

We utilized two schemes for heavy quarks treatment:

- FFNS - with only light quarks active
- VFNS - with the number of active quarks depending on the process energy scale (here it doesn't change as the total cross-sections are always computed in FFNS and $p_T = const$ in case of differential cross-sections)

We used two loop formula for α_s and the WW function of Equivalent Photons Approximation.

Each total cross-section was calculated with various values of the process energy scale μ and various values of masses of the heavy quarks. Also we did our computation with and without a cut on the part of electron energy carried by the photon y .

FFNS

HERA

μ^2 :	$1/4(p_T^2 + m_c^2)$	$p_T^2 + m_c^2$	$4(p_T^2 + m_c^2)$
$m_c = 1.2$ GeV	632	825	939
$m_c = 1.5$ GeV	602	694	732
$m_c = 1.8$ GeV	386	425	438



Table 1: HERA: $E_p = 920$ GeV, $E_e = 27.5$ GeV. All values in nb. LO calculated $\sigma(ep \rightarrow cc/\bar{c}X)$. GRVHO and MRS(G)HO parametrizations of the real photon and proton were used. $0 < y < 1$

THERA

70

(c)

μ^2 :	$1/4(p_T^2 + m_c^2)$	$p_T^2 + m_c^2$	$4(p_T^2 + m_c^2)$
$m_c = 1.2$ GeV	2956	4315	5368
$m_c = 1.5$ GeV	1571	2126	2532
$m_c = 1.8$ GeV	899	1146	1317

nb

Table 2: TERA: $E_p = 920$ GeV, $E_e = 250$ GeV. All values in nb. LO calculated $\sigma(ep \rightarrow ec/\bar{c}X)$. GRVHO and MRS(G)HO parametrizations of the real photon and proton were used. $0 < y < 1$

μ^2 :	$1/4(p_T^2 + m_c^2)$	$p_T^2 + m_c^2$	$4(p_T^2 + m_c^2)$
$m_c = 1.2$ GeV	866	1423	1928
$m_c = 1.5$ GeV	505	749	957
$m_c = 1.8$ GeV	315	432	526

Table 3: TERA: $E_p = 920$ GeV, $E_e = 250$ GeV. All values in nb. LO calculated $\sigma(ep \rightarrow ec/\bar{c}X)$. GRVHO and MRS(G)HO parametrizations of the real photon and proton were used. $0.15 < y < 0.86$

nb

(b)

μ^2 :	$1/4(p_T^2 + m_b^2)$	$p_T^2 + m_b^2$	$4(p_T^2 + m_b^2)$
$m_b = 4.25$ GeV	30.6	30.5	29.8
$m_b = 4.5$ GeV	26.1	25.9	25.2
$m_b = 4.75$ GeV	22.4	22.1	21.5

nb

Table 4: TERA: $E_p = 920$ GeV, $E_e = 250$ GeV. All values in nb. LO calculated $\sigma(ep \rightarrow eb/\bar{b}X)$. GRVHO and MRS(G)HO parametrizations of the real photon and proton were used. $0 < y < 1$

$\mu^2:$	$1/4(p_T^2 + m_b^2)$	$p_T^2 + m_b^2$	$4(p_T^2 + m_b^2)$
$m_b = 4.25$ GeV	14.1	14.6	14.7
$m_b = 4.5$ GeV	12.1	12.5	12.6
$m_b = 4.75$ GeV	10.5	10.8	10.8

Table 5: TERA: $E_p = 920$ GeV, $E_e = 250$ GeV. All values in nb. LO calculated $\sigma(ep \rightarrow eb/\bar{b}X)$. GRVHO and MRS(G)HO parametrizations of the real photon and proton were used. $0.15 < y < 0.86$



$\mu^2:$	$1/4(p_T^2 + m_t^2)$	$p_T^2 + m_t^2$	$4(p_T^2 + m_t^2)$
$m_t = 170$ GeV	0.01264	0.00995	0.00801
$m_t = 174$ GeV	0.01041	0.00818	0.00657
$m_t = 178$ GeV	0.00858	0.00672	0.00539

pb

Table 6: TERA: $E_p = 920$ GeV, $E_e = 250$ GeV. All values in pb. LO calculated $\sigma(ep \rightarrow et/\bar{t}X)$. GRVHO and MRS(G)HO parametrizations of the real photon and proton were used. $0 < y < 1$

$\mu^2:$	$1/4(p_T^2 + m_t^2)$	$p_T^2 + m_t^2$	$4(p_T^2 + m_t^2)$
$m_t = 170$ GeV	0.00880	0.00686	0.00548
$m_t = 174$ GeV	0.00715	0.00556	0.00443
$m_t = 178$ GeV	0.00581	0.00451	0.00359

Table 7: TERA: $E_p = 920$ GeV, $E_e = 250$ GeV. All values in pb. LO calculated $\sigma(ep \rightarrow et/\bar{t}X)$. GRVHO and MRS(G)HO parametrizations of the real photon and proton were used. $0.15 < y < 0.86$

direct / resolved

P. Jankowski

$ep \rightarrow c/\bar{c} \chi$ THERA

$p_T = 10 \text{ GeV}$

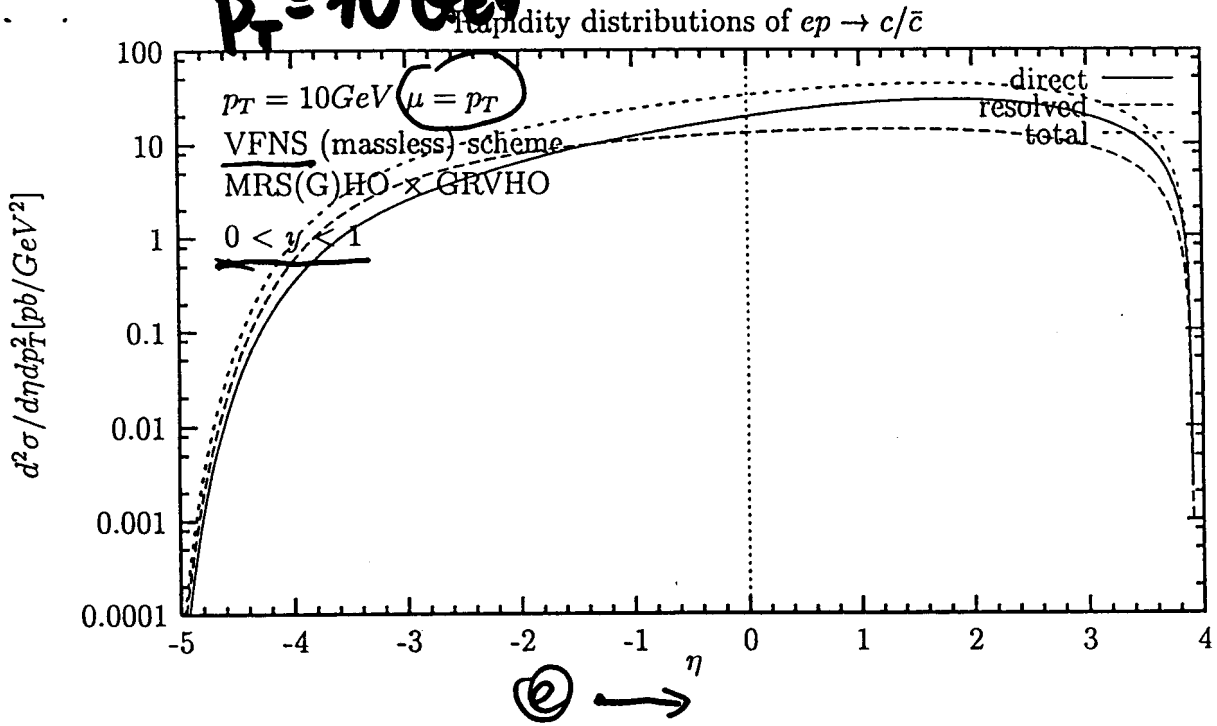


Figure 3: c/\bar{c} production, TERA, $E_p = 920 \text{ GeV}$, $E_e = 250 \text{ GeV}$

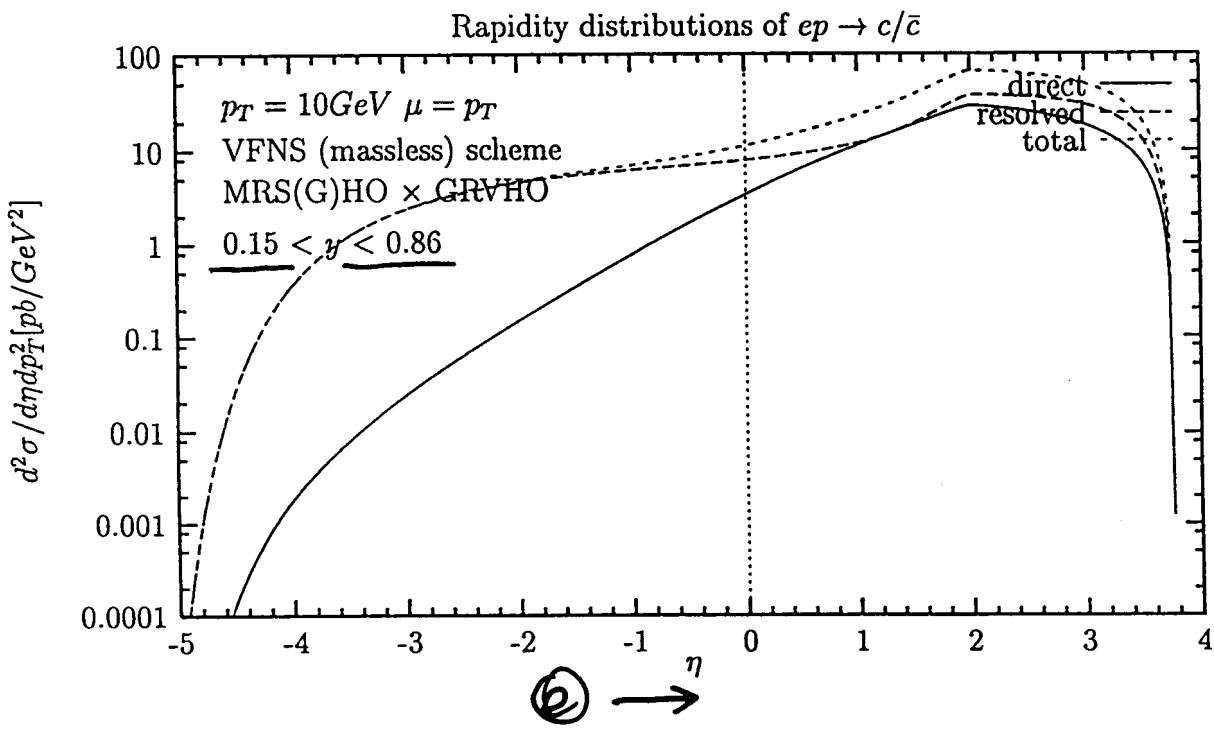


Figure 4: c/\bar{c} production, TERA, $E_p = 920 \text{ GeV}$, $E_e = 250 \text{ GeV}$

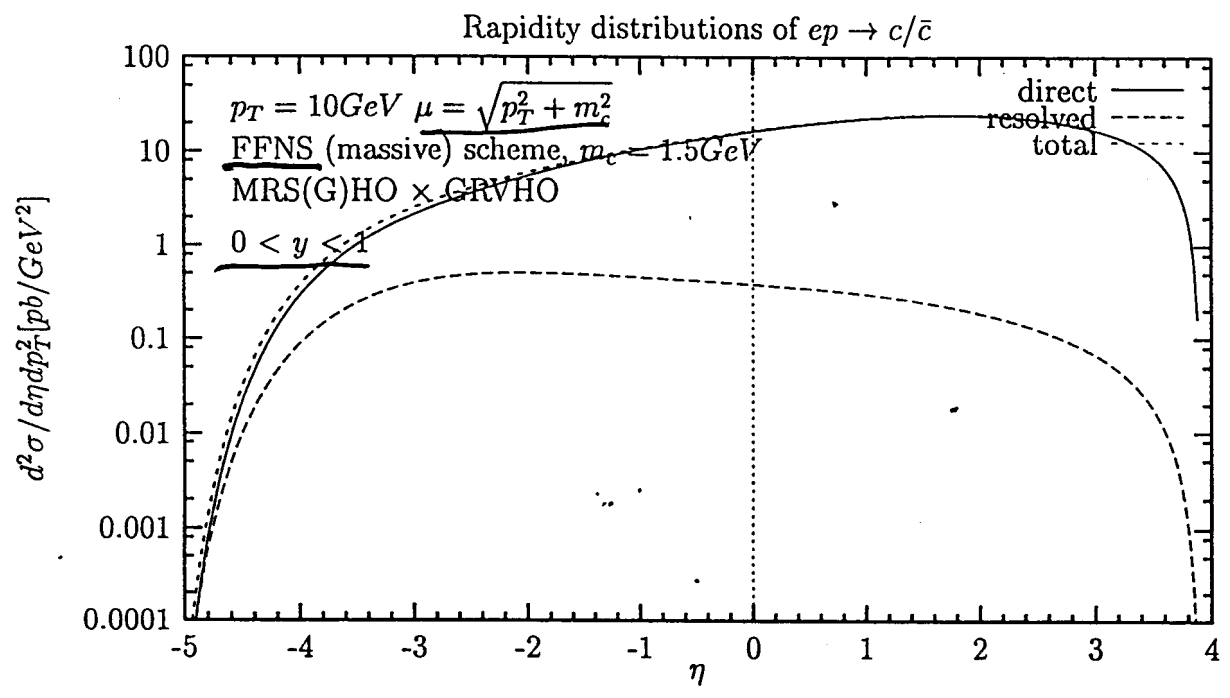


Figure 1: c/\bar{c} production, TERA, $E_p = 920 \text{ GeV}$, $E_e = 250 \text{ GeV}$

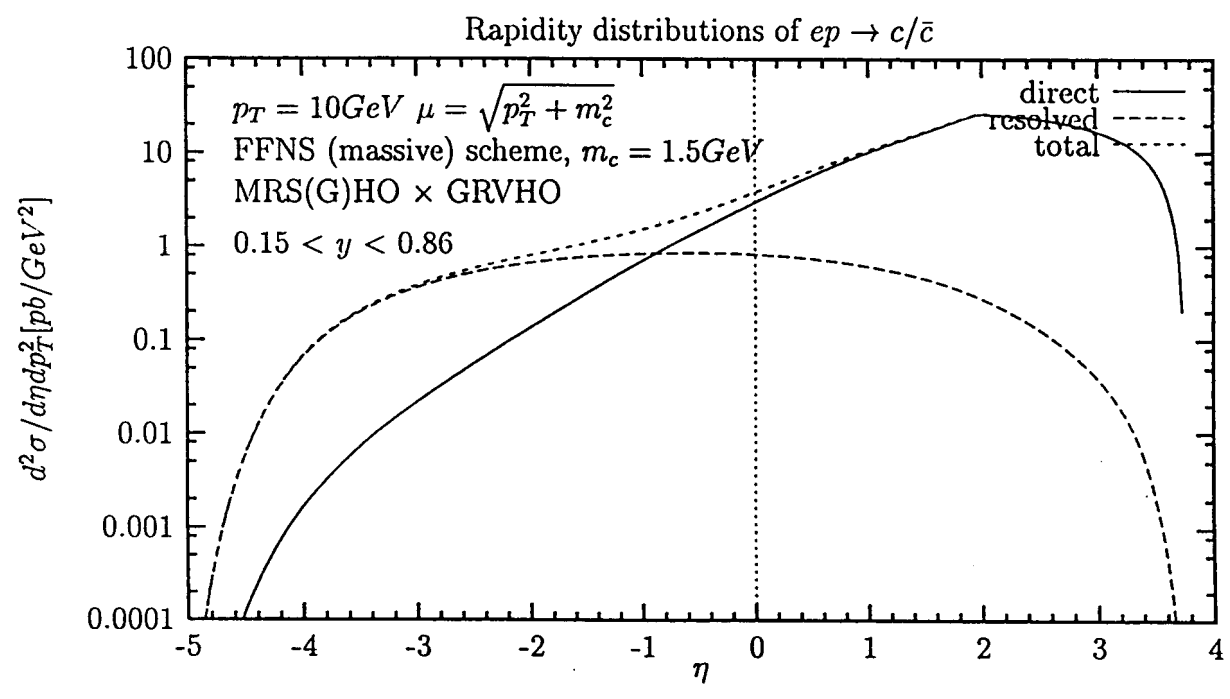


Figure 2: c/\bar{c} production, TERA, $E_p = 920 \text{ GeV}$, $E_e = 250 \text{ GeV}$

b

$p_T = 10 \text{ GeV}$

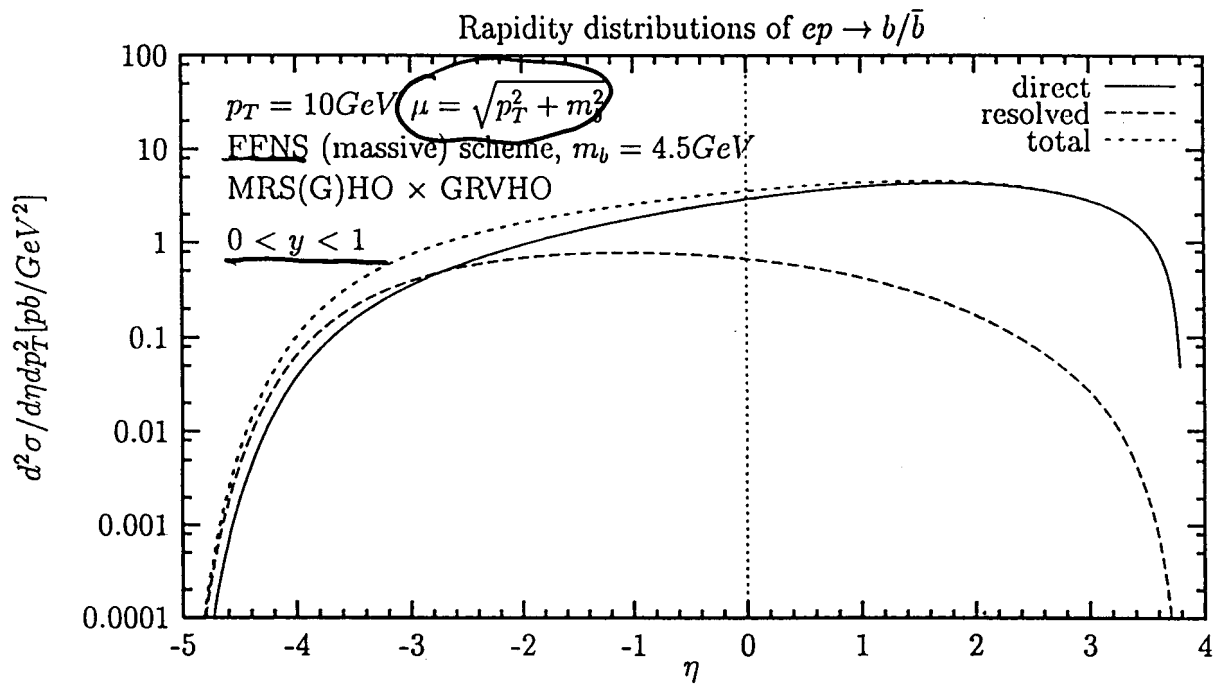


Figure 5: b/\bar{b} production, TERA, $E_p = 920 \text{ GeV}$, $E_e = 250 \text{ GeV}$

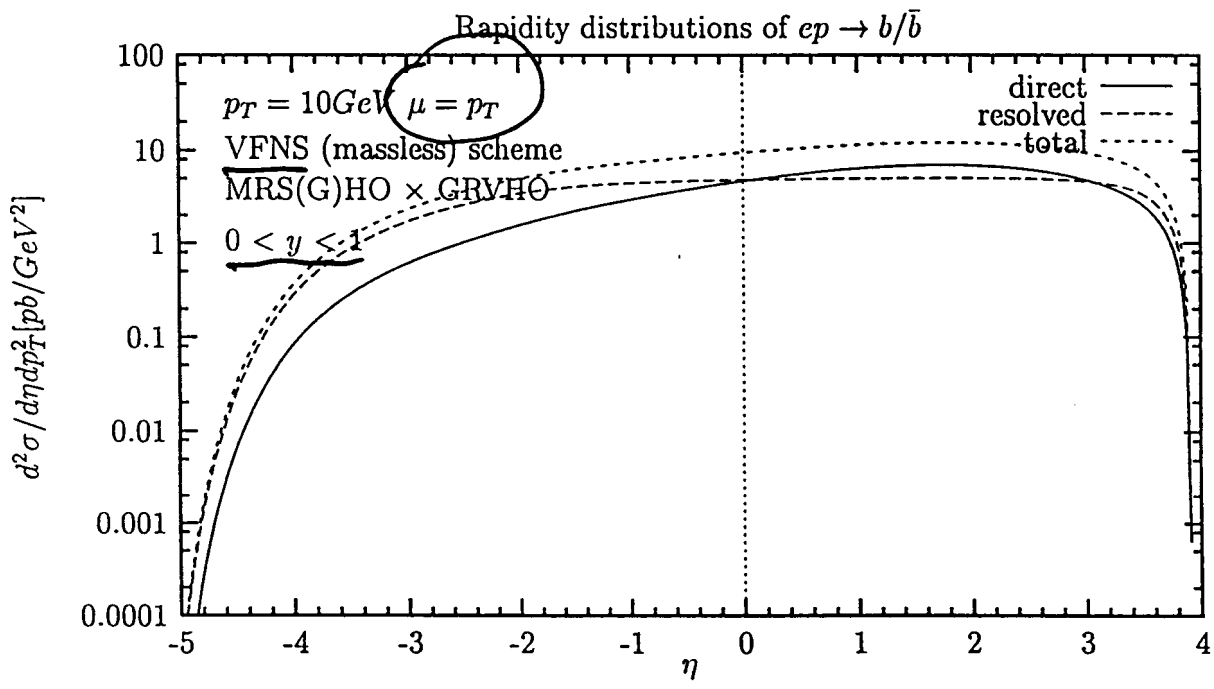


Figure 6: b/\bar{b} production, TERA, $E_p = 920 \text{ GeV}$, $E_e = 250 \text{ GeV}$

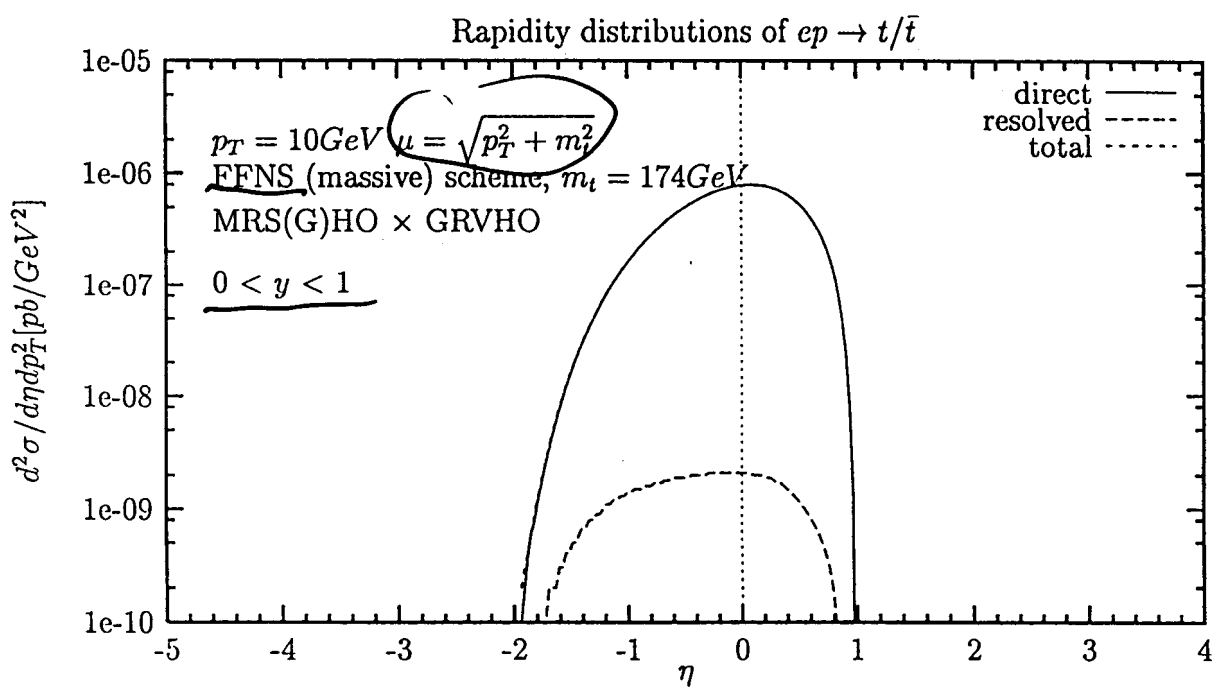
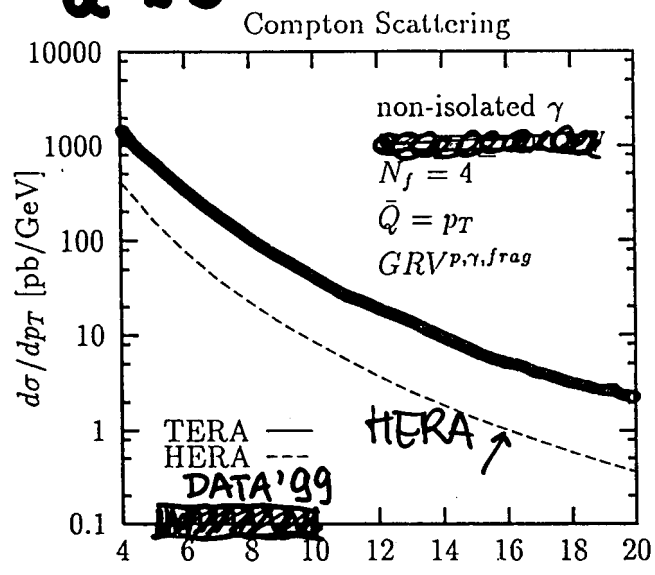
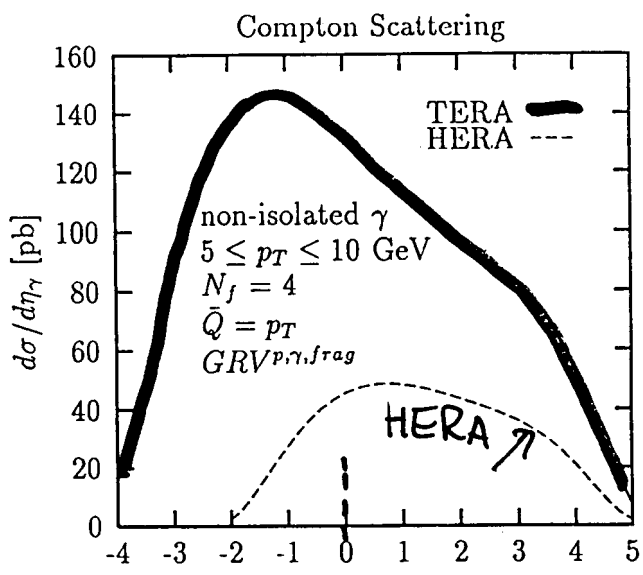


Figure 7: t/\bar{t} production, TERA, $E_p = 920 \text{ GeV}$, $E_e = 250 \text{ GeV}$

Prompt photons at THERA

- photoproduction $Q^2 \approx 0$

A. Zembriusl
 NLO calc. ^{MK}



$p \rightarrow \eta \gamma$

p_T, GeV

- also for $Q^2 \neq 0$ ($Q^2 \ll p_T^2$)

↳ virtual photon γ^*
 with virtuality $Q^2 (= -Q^2)$
 up to few GeV^2

Can we probe individual parton density
 in γ ? in γ^* ?

↳

Prompt photon production (Compton) process

A. Zembruski
MK

LO calc.
(hep-ph/9312368)
PRD 57(98)

virtuality of γ
 $P^2 \approx 0 - 1.0$
 GeV^2

Ratio: $\frac{q_\gamma^\sigma q_p^P \rightarrow \gamma q_\gamma}{\gamma q_\gamma^P \rightarrow \gamma q_\gamma}$ ← Born

and for

- $q_\gamma^\sigma q_p^P \rightarrow \gamma q_\gamma$
- $\bar{q}_\gamma^\sigma q_p^P + q_\gamma^\sigma \bar{q}_p^P \rightarrow \gamma q_\gamma$

we study the diff. cross sections

$\frac{d\sigma}{dE_\gamma dY dP^2 d\eta_\gamma}$

$q_\gamma^\sigma q_p^P \rightarrow \gamma q_\gamma$
dominate⁵
here

HERA

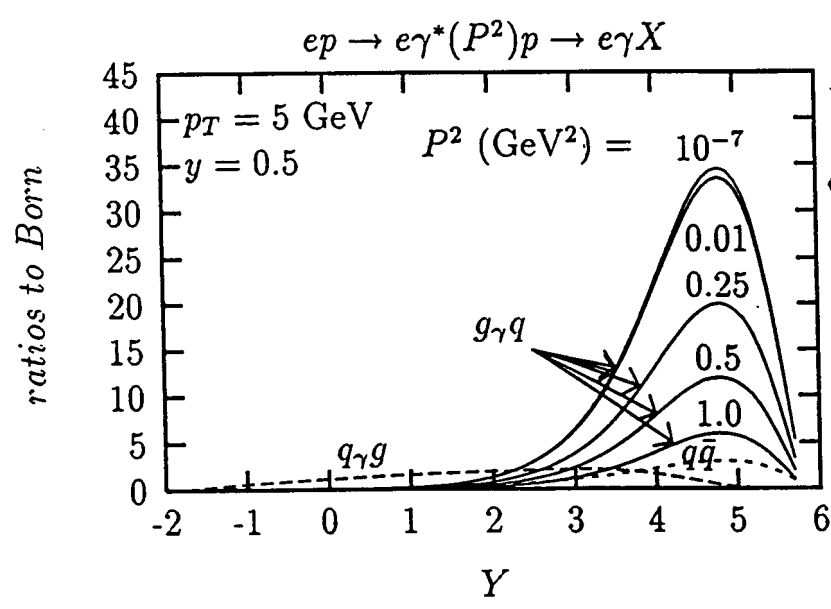
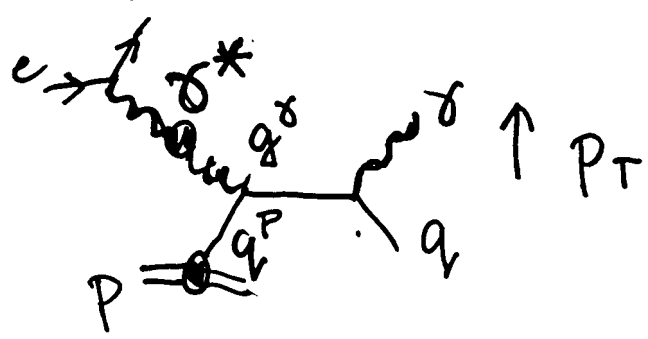


Figure 2. Ratios of the cross sections (see eq. 13) for the reaction $ep \rightarrow e\gamma X$, with $y=0.5$ and $p_T = 5 \text{ GeV}$. The ratios of the contributions due to $g_\gamma \cdot q_p \rightarrow \gamma q$ (solid lines), $q_\gamma \cdot g_p \rightarrow \gamma q$ (long-dashed line) and $q_\gamma \cdot \bar{q}_p \rightarrow \gamma q$ (short-dashed line) to the Born cross section are presented for various values of the virtuality P^2 as a function of the rapidity Y . For $g_\gamma \cdot q_p \rightarrow \gamma q$ the results correspond to $P^2 = 10^{-7}, 0.01, 0.25, 0.5$ and 1 GeV^2 . For $q_\gamma \cdot g_p \rightarrow \gamma q$ and $q_\gamma \cdot \bar{q}_p \rightarrow \gamma q$ we take $P^2 = 0.1 \text{ GeV}^2$.

$ep \rightarrow \gamma X$

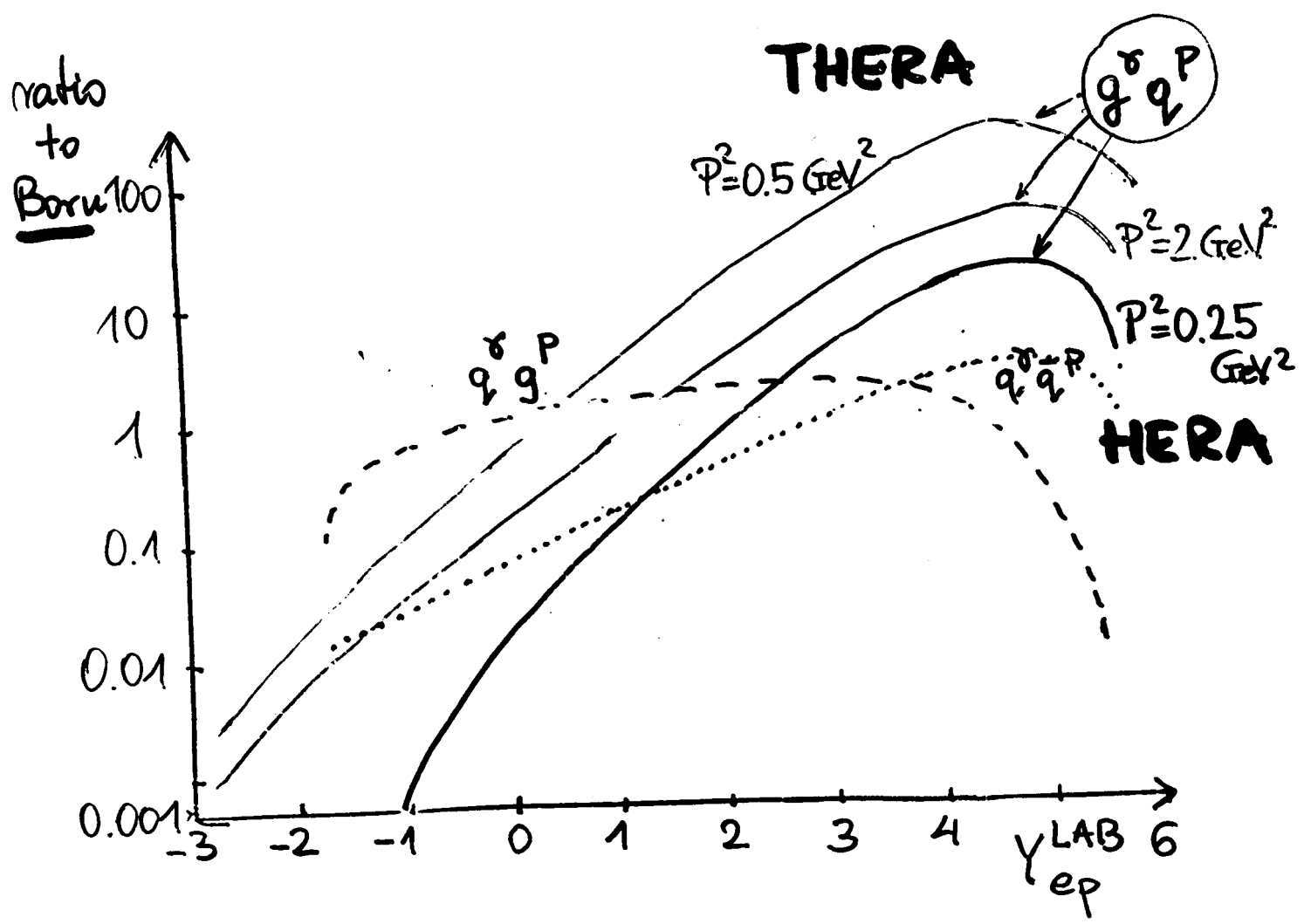


γ^* with virtuality P^2

$P_T = 5 \text{ GeV}$

$y = 0.5$

HERA \rightarrow THERA

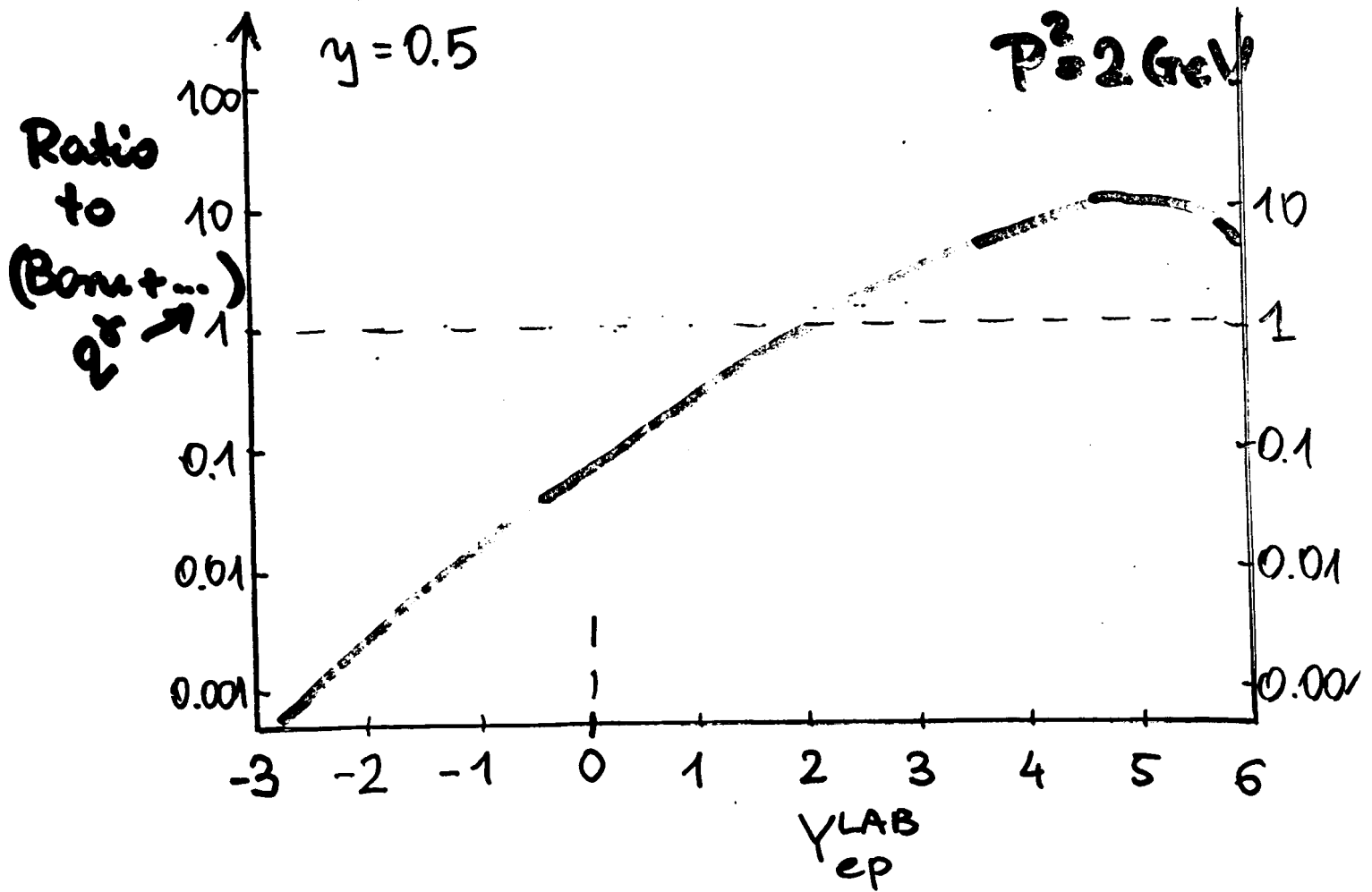


THERA

$$P_T = 5 \text{ GeV}$$

$$y = 0.5$$

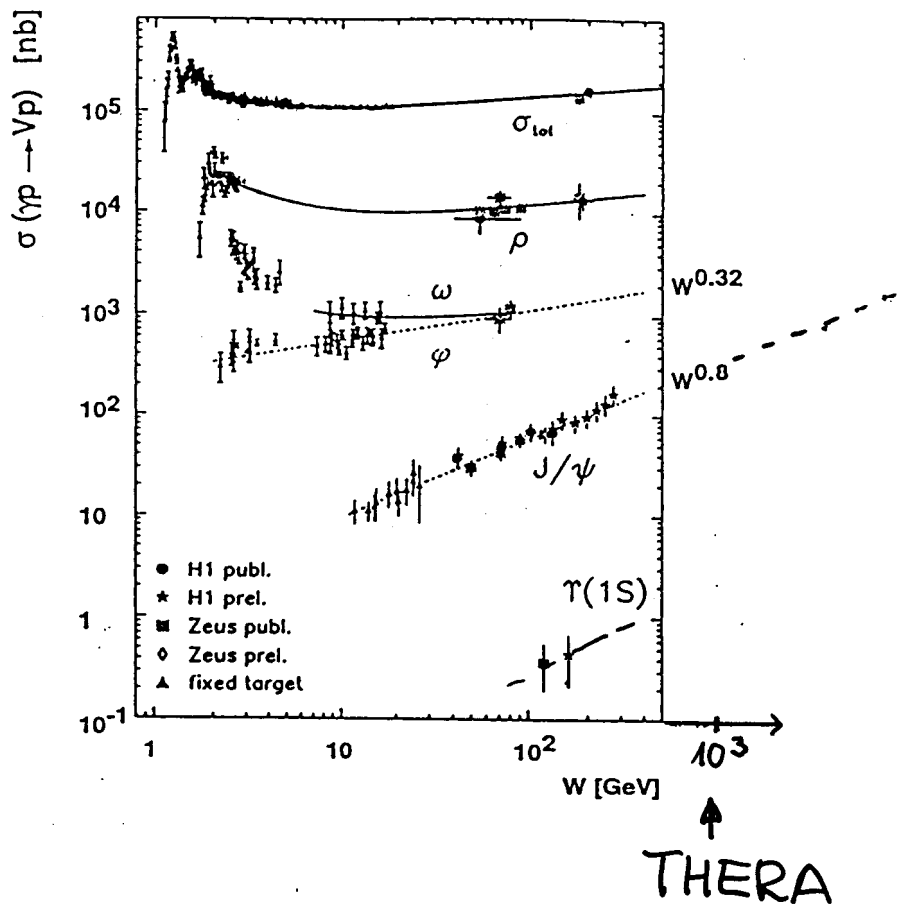
$$P^2 = 2 \text{ GeV}^2$$



Exclusive photoproduction

$$\gamma p \rightarrow V p$$

$$V = \rho, \omega, \varphi, J/\psi, \Upsilon$$



$$\gamma p \rightarrow \pi^0 p$$

$$\uparrow \eta, \eta', \eta_c \dots$$

→ J. Ginzburg talk