Pentagon diagrams give terms

$$\frac{1}{(d-4)} I_{11111}^{(d)}$$

The reason is terms in integrand like

$$\int d^d k_1 \frac{(\overline{U}\ldots\hat{k}_1\ldots U)(\overline{V}\ldots\hat{k}_1\ldots V)}{c_1c_2c_3c_4c_5} \qquad (1)$$

$$\begin{aligned}\int d^d k_1 \frac{k_{1\mu}k_{1\nu}}{c_1c_2c_3c_4c_5} = -\frac{1}{2} g_{\mu\nu} I_{11111}^{(d+2)} \\ + 2p_{1\mu}p_{1\nu} I_{31111}^{(d+4)} + 2I_{13111}^{(d+4)} p_{2\mu}p_{2\nu} \\ + 2I_{11311}^{(d+4)} p_{3\mu}p_{3\nu} + 2I_{11131}^{(d+4)} p_{4\mu}p_{4\nu} \\ + I_{22111}^{(d+4)}(p_{1\mu}p_{2\nu} + p_{1\nu}p_{2\mu}) + I_{21211}^{(d+4)}(p_{1\mu}p_{3\nu} + p_{1\nu}p_{3\mu}) \\ + I_{21121}^{(d+4)}(p_{1\mu}p_{4\nu} + p_{1\nu}p_{4\mu}) + I_{12211}^{(d+4)}(p_{2\mu}p_{3\nu} + p_{2\nu}p_{3\mu}) \\ + I_{12121}^{(d+4)}(p_{2\mu}p_{4\nu} + p_{2\nu}p_{4\mu}) + I_{11221}^{(d+4)}(p_{3\mu}p_{4\nu} + p_{3\nu}p_{4\mu})\end{aligned}$$

By using recurrence relations one gets

$$I_{11111}^{(d+2)} = \frac{2\Delta_5}{(d-4)G_4} I_{11111}^{(d)} + \sum_{k=1}^{5} \frac{\partial_k \Delta_5}{(d-4)G_4} \mathbf{k}^- I_5^{(d)}$$

There were no such problems for 3,4- point functions. Multiplying (1) by Born terms, summing over polarizations, taking traces removes problematic $1/(d-4)$ terms.

Another possibility will be to use basis of master integrals in higher dimensions