

Measuring the Beam Energy with Radiative Return Events

Master Thesis of Arnd Hinze
(presented by Klaus Moenig at LCWS05)

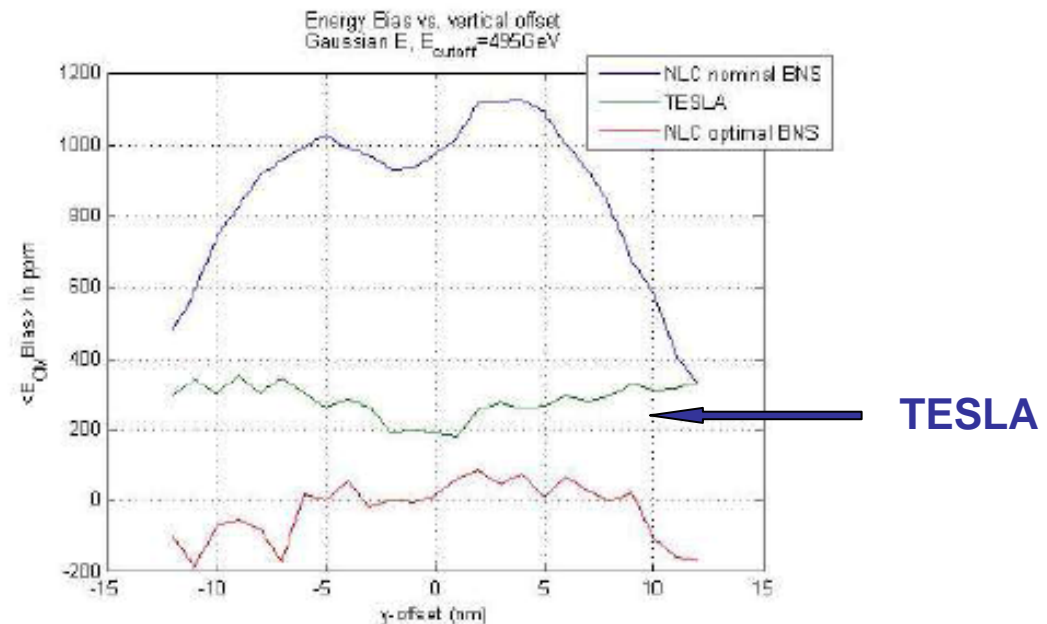
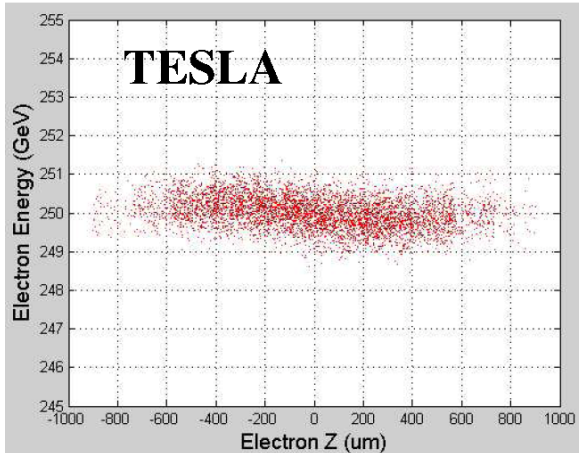
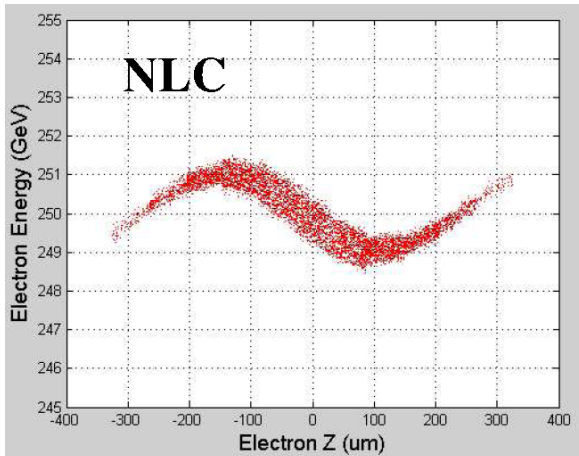
Introduction

- The beam energy is needed on the 10^{-4} level for mass determinations at the ILC (Higgs, top, SUSY...)
- The beam-energy measurement will mainly come from a magnetic spectrometer
- The absolute calibration of the spectrometer is difficult
- The sum of the beam energies is not necessarily equal to the luminosity weighted centre of mass energy \sqrt{s} → next slide
- It would thus be useful to have a method to determine the cms energy from real (annihilation) data
- Such a method exists: radiative return events $e^+e^- \rightarrow Z \gamma \rightarrow f \bar{f} \gamma$
- The validity of this method was already proven at LEP

Energy Bias from Kink-Instabilities

A. Florimonte, M. Woods, IPBI TN-2005-1

- Wakefields introduce a correlation $z - E$
 - Disruption give a different weight to different parts of the bunch
- These effects make the luminosity weighted cms energy different from twice the beam energy



500 GeV TESLA: mean = 150 ppm
 spread = 30 ppm
 max = 350 ppm

Basic Idea of the Radiative Return Analysis

- The Z-mass is known with very high precision from LEP
- assume only one photon is radiated

⇒ $\sqrt{s'}$ can be calculated from fermion angles only

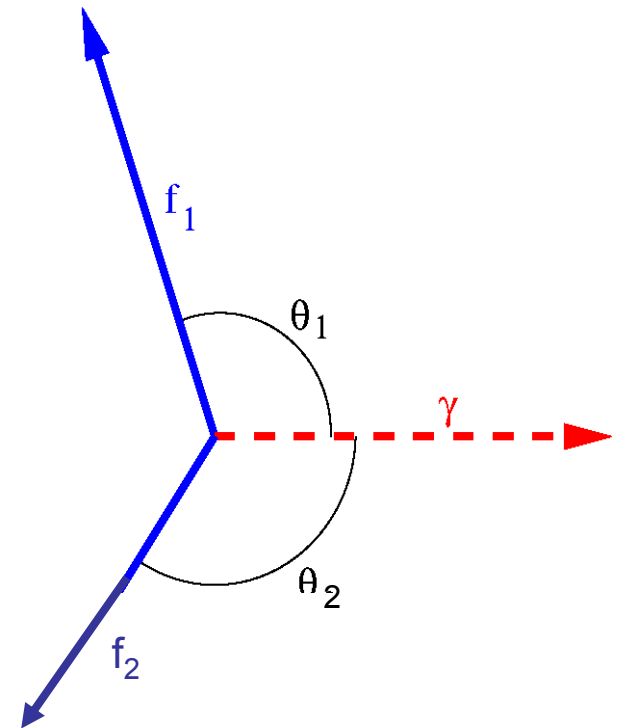
$$\frac{\sqrt{s'}}{\sqrt{s}} = \sqrt{\frac{\sin \theta_1 + \sin \theta_2 + \sin(\theta_1 + \theta_2)}{\sin \theta_1 + \sin \theta_2 - \sin(\theta_1 + \theta_2)}}$$

- γ either along the beampipe or its angle is measured
- this formula assumes that the fermion mass can be neglected

- Assume $\sqrt{s'} = M_Z$

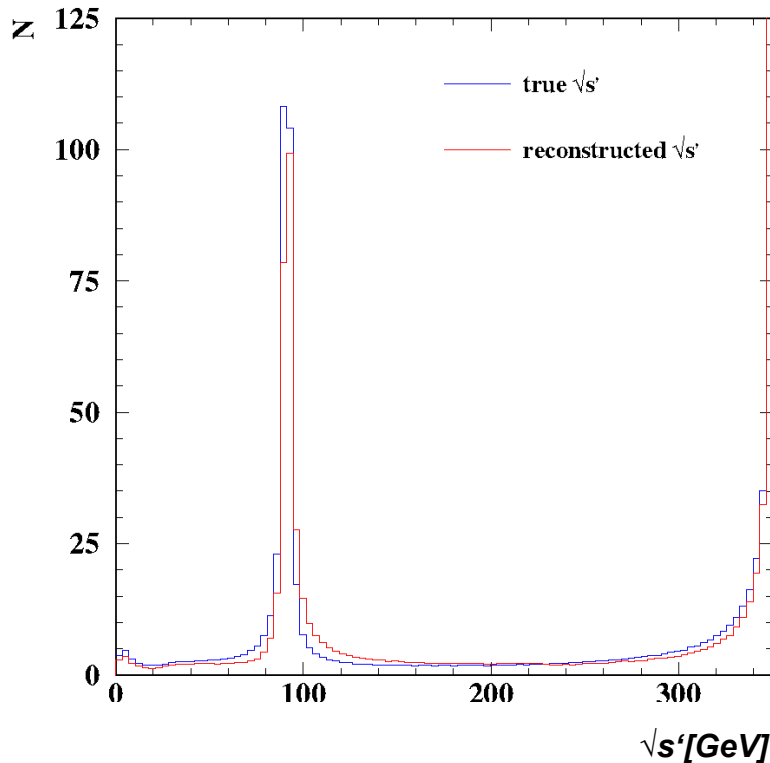


$$\sqrt{s} = M_Z \sqrt{\frac{\sin \theta_1 + \sin \theta_2 + \sin(\theta_1 + \theta_2)}{\sin \theta_1 + \sin \theta_2 - \sin(\theta_1 + \theta_2)}}$$

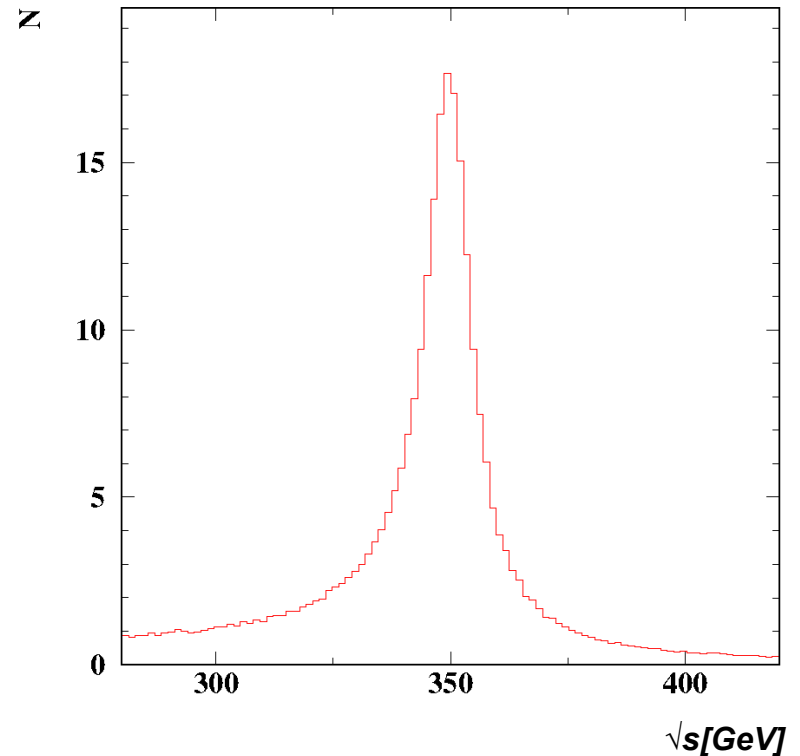


First analysis with $e^+e^- \rightarrow \mu^+\mu^-\gamma$ at $\sqrt{s} = 350$ GeV and int. lumi = 100 fb^{-1}

$\sqrt{s'}$ spectrum



Reconstructed \sqrt{s}



- $\sigma(\text{rad. ret.}) \sim 0.5 \text{ pb}$, scales approx. with $1/s$
- detector efficiency (7° cut) $\approx 90\%$
- ideal beam with beamstrahlung and 0.2% Gaussian energy spread

Method used: Fit Procedure

- χ^2 fit of 'data' to MC (reference) samples with free normalisation
 - 'data' sample at $\sqrt{s}_{\text{dat}} = 350$ GeV, ***MC_{dat}***
 - and reference sample at $\sqrt{s}_{\text{ref}} = \sqrt{s}_{\text{dat}} + \Delta\sqrt{s}$ (= 1GeV), ***MC_{ref}***

$$N_{\text{pred}}(\sqrt{s}) = \frac{\sqrt{s_{\text{ref}}} - \sqrt{s_{\text{dat}}}}{\Delta\sqrt{s}} (MC_{\text{ref}} - MC_{\text{dat}})$$

- easy to include all effects into fit
- fit tested to be bias free in region $\sqrt{s}_{\text{dat}} \pm \Delta\sqrt{s}$

Cuts: • $7^\circ < \theta_{1,2} < 183^\circ$ (detector acceptance of muons)

- $M_Z - 5 \text{ GeV} < m(\mu^+\mu^-) < M_Z + 5 \text{ GeV}$

Backgrounds

$\gamma\gamma$ background ($e^+e^- \rightarrow e^+e^-\mu^+\mu^-$):

- very large cross section without cuts
- less than 10% background in fit range after $m(\mu^+\mu^-)$ cut

Zee:

- cross section similar to signal
- after cut on visible electrons $\sim 25\%$ background remains
- however kinematics similar to signal, so no problem

WW, ZZ:

- ZZ already small, WW reduced by $m(\mu^+\mu^-)$ cut
- in the end $\sim 1\%$ background

Results

Fit to 100 fb^{-1} including beam effects and background:

**stat.
error**

$$\Delta\sqrt{s} = 47 \text{ MeV} \quad \left(\frac{\Delta\sqrt{s}}{\sqrt{s}} = 1.3 \cdot 10^{-4} \right)$$

$$\sqrt{s} = 350 \text{ GeV}$$

- without beamstrahlung and energy spread ~10% better
- little effect from background
- slight improvement possible if 2D fit (\sqrt{s} , θ)

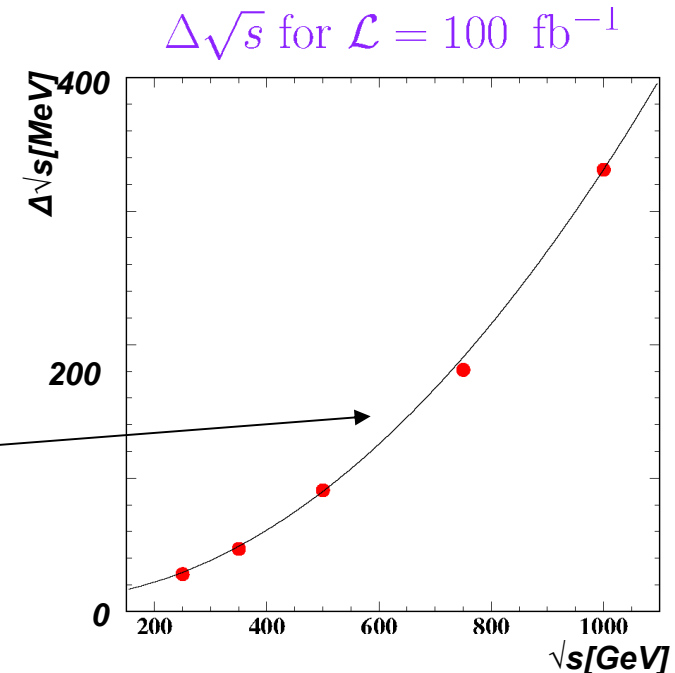
Result strongly energy dependent:

- cross section proportional $1/s$
- resolution deteriorates with larger s
- background rises with s
- acceptance worse at larger s

Parameterisation:

$$\Delta\sqrt{s} = (8.8 + 0.0026\sqrt{s}/\text{GeV} + 0.0032s/\text{GeV}^2) \text{ MeV}$$

H.Juergen Schreiber DESY
23/05/05



Systematics

Background: no effect for 20-30% background uncertainty

Energy spread: $\Delta\sqrt{s} = 10$ MeV if Gaussian energy spread is replaced by rectangular, no effect if 0.1% instead of 0.2%

Beamstrahlung: method largely cancels errors from beamstrahlung determination

Aspect ratio of tracker: LEP error $\Delta(\delta R/\delta L) = \delta \tan\theta = 5 \cdot 10^{-4}$



$\Delta\sqrt{s} = 160$ MeV, which has to be about an order of magnitude better !

Conclusions

The centre of mass energy can be measured on the 10^{-4} level from radiative return events.

Inclusion of Bhabha scattering $e^+e^- \rightarrow e^+e^-\gamma$, $e^+e^- \rightarrow \tau^+\tau^-\gamma$ or $e^+e^- \rightarrow q \bar{q} \gamma$ might improve the result.

But each of this channel contains additional complications, which dilute possible improvements.

This is a high luminosity analysis, so relative measurements e.g. in a scan are needed from spectrometers.

The length to radius ratio of the tracking detector needs to be known to better than 10^{-4} not to be limited by this effect.

A global analysis of Bhabha acolinearity for beamstrahlung and radiative return events for the beam energy is needed to understand effects from beam-beam correlations.