

Two-dimensional harmonic polylogarithms for Bhabha scattering

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Outline

Methods of calculations for Feynman integrals in Particle Physics

Numerical and analytical methods

Differential equation method

Harmonic and Generalised Harmonic polylogarithms

Definitions of HPLs

Definitions of GPLs

GPLs for the Bhabha process

Final remarks

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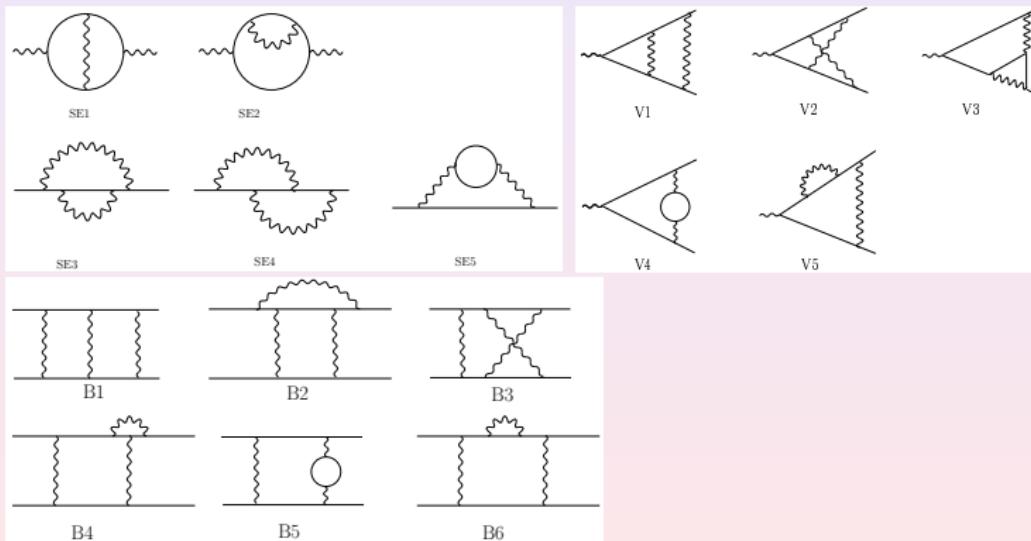
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Bhabha scattering $e^+e^- \rightarrow e^+e^-$, $d = 4 - 2\epsilon$

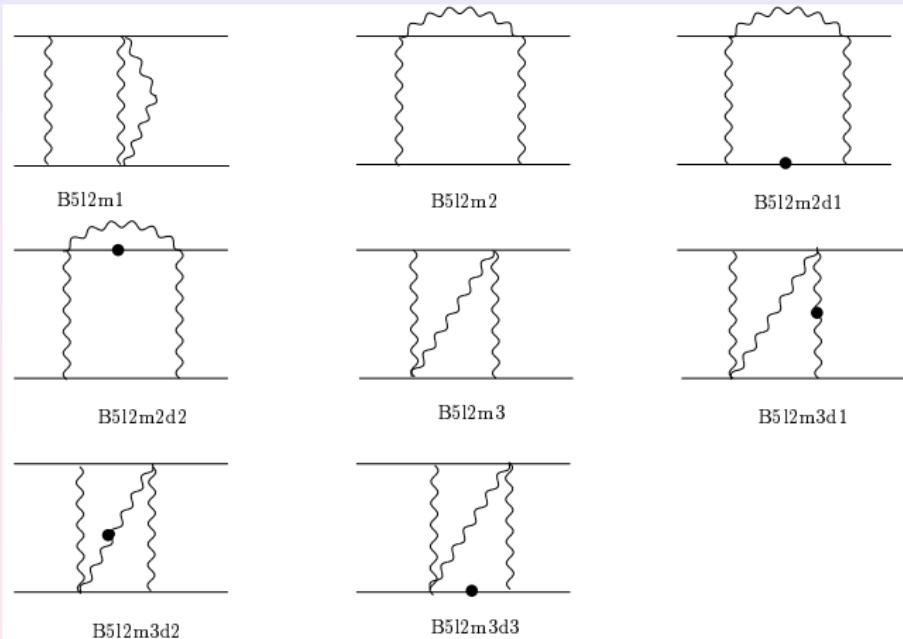
$$G(X) = \frac{1}{(i\pi^{d/2})^L} \int \frac{d^d k_1 \dots d^d k_L \ X}{(q_1^2 - m_1^2)^{\nu_1} \dots (q_j^2 - m_j^2)^{\nu_j} (q_N^2 - m_N^2)^{\nu_N}},$$



Bhabha scattering $e^+e^- \rightarrow e^+e^-$, Czakon, JG, Riemann, PRD2005

MI	B1	B2	B3	B4	B5	B6	Solved
B714m1	+	-	-	-	-	-	Y
B714m1N	+	-	-	-	-	-	Y
B714m2	-	+	-	-	-	-	Y/N
B714m2[d1--d3]	-	+	-	-	-	-	Y/N
B714m3	-	-	+	-	-	-	Y/N
B714m3[d1--d2]	-	-	+	-	-	-	
B613m1	+	-	+	-	-	-	
B613m1d	+	-	+	-	-	-	
B613m2	-	+	-	+	-	-	
B613m2d	-	+	-	+	-	-	
B613m3	-	-	+	-	-	-	
B613m3[d1--d5]	-	-	+	-	-	-	
B512m1	+	-	+	-	-	-	Y
B512m2	-	+	-	+	-	+	Y/N
B512m2[d1--d2]	-	+	-	+	-	+	Y/N
B512m3	+	-	+	-	-	-	Y/N
B512m3[d1--d3]	+	-	+	-	-	-	Y/N
B513m	-	+	+	+	-	-	
B513m[d1--d3]	-	+	+	+	-	-	
B514m	-	+	+	+	+	-	Y
B514md	-	+	+	+	+	-	Y
B412m*	-	-	-	+	+	+	Y

Bhabha scattering $e^+e^- \rightarrow e^+e^-$, Czakon, JG, Riemann, PRD2005



Numerical Methods

- ▶ **No problem** with the presence of different scales
- ▶ **Problem:** infrared singularities, thresholds and pseudothresholds

Solved:

- ▶ two-loop SE and vertex diagrams, multi-leg one-loop diagrams
Ghinculov et al, 1995, *Passarino et al*, 2001
- ▶ Sector decomposition method: multiloop calculations in the non-physical region and evaluation of phase-space integrals
Binoth-Heinrich, 2000
- ▶ SE by DEqs with Runge-Kutta method
Caffo et al, 2002

Analytical Methods

- ▶ **Problem:** At most three different scales
- ▶ **No problem** with special regions

New methods:

- ▶ *Laporta-Remiddi* method: Master Integrals
- ▶ Expansion by regions
V. Smirnov
- ▶ Mellin-Barnes transformation
Davydychev, Smirnov
- ▶ Differential Equation Method
Kotikov, Remiddi

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General idea

M with respect to the external scale x ,

$$\frac{d}{dx}M(x) = A(x)M(x) + B(x).$$

If $H(x)$ is a solution of the homogeneous equation

$$\frac{d}{dx}H(x) = A(x)H(x).$$

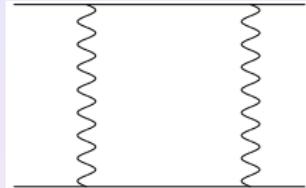
then the full solution is given by

$$M(x) = H(x) \left(\int^x H^{-1}(x')B(x')dx' + \text{Const} \right)$$

Solutions are generally combinations of [general] hypergeometric functions which are difficult to expand in powers of ϵ

1-loop box,

Fleischer, Jegerlehner, Tarasov



$$\begin{aligned}
 &\sim \frac{t - 4m^2}{i\pi^{d/2}\Gamma\left(2 - \frac{d}{2}\right)} \int \frac{d^d k_1}{k^2(k^2 + 2q_4 k)(k + q_1 + q_4)^2(k^2 - 2q_3 k)} \\
 &= \dots \frac{\Gamma\left(\frac{d-2}{2}\right)}{\Gamma\left(\frac{d-1}{2}\right)} F_1\left(\frac{d-3}{2}, 1, \frac{1}{2}; \frac{d-1}{2}; \frac{sZ(s, u, t)}{4}, 1 - \frac{s}{4m^2}\right) \\
 &\quad + \dots F_2\left(\frac{d-3}{2}, 1, 1, \frac{3}{2}, \frac{d-2}{2}; \frac{t}{t - 4m^2}, -m^2 Z(s, u, t)\right) \\
 &\quad + \dots F_{1;1;0}^{1;2;1} \left[\frac{\frac{d-3}{2}; \frac{d-3}{2}}{\frac{d-1}{2}; \frac{d-2}{2}}; 1; 1; -; -m^2 Z(s, u, t), 1 - \frac{4m^2}{s} \right]
 \end{aligned}$$

ϵ expansion:

Fleischer, Riemann, Tarasov, APP. B34, 5345

Appell hypergeometric functions:

$$F_1 \left(\frac{d-3}{2}, 1, \frac{1}{2}, \frac{d-1}{2}; x, y \right) = \sum_{r=0}^{\infty} \sum_{s=0}^{\infty} \frac{\left(\frac{d-3}{2}\right)_{r+s}}{\left(\frac{d-1}{2}\right)_{r+s}} \frac{\left(\frac{1}{2}\right)_s}{(1)_s} x^r y^s,$$

$$F_2 \left(\frac{d-3}{2}, 1, 1, \frac{3}{2}, \frac{d-2}{2}; x, y \right) = \sum_{r=0}^{\infty} \sum_{s=0}^{\infty} \frac{\left(\frac{d-3}{2}\right)_{r+s}}{\left(\frac{3}{2}\right)_r \left(\frac{d-2}{2}\right)_s} x^r y^s$$

the Kampé de Fériet function:

$$F_{1;1;0}^{1;2;1} \left[\begin{matrix} \frac{d-3}{2}; & \frac{d-3}{2}, 1; & 1; \\ \frac{d-1}{2}; & \frac{d-2}{2}; & -; \end{matrix} x, y \right] = \sum_{r=0}^{\infty} \sum_{s=0}^{\infty} \frac{\left(\frac{d-3}{2}\right)_{r+s}}{\left(\frac{d-1}{2}\right)_{r+s}} \frac{\left(\frac{d-3}{2}\right)_r}{\left(\frac{d-2}{2}\right)_r} x^r y^s.$$

Order by order expansion in ϵ

$$M(x) = H(x) \left(\int^x H^{-1}(x') B(x') dx' + \text{Const} \right)$$

$$M = \sum_{i=-m}^n m^i \epsilon^i, \quad A = \sum_{i=0}^{n+\alpha} a^i \epsilon^i, \quad B = \sum_{i=-\alpha}^n b^i \epsilon^i$$

then

$$\boxed{\frac{d}{dx} m^i(x) = \sum_{j=0}^{m-i} a^j(x) m^{i-j}(x) + b^i(x)}$$

Conformal transformation

In physical situation, the differential operator is:

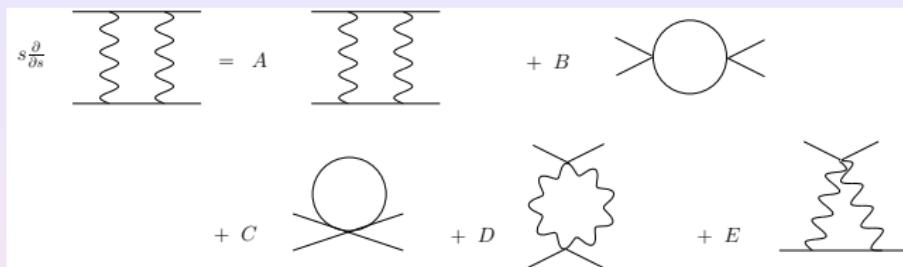
$$s \frac{\partial}{\partial s} = \frac{1}{2} \left\{ (p_1^\mu + p_2^\mu) + \frac{s(p_2^\mu - p_3^\mu)}{s+t-4} \right\} \frac{\partial}{\partial p_2^\mu}.$$

Transformation

$$\begin{aligned}x &= \frac{\sqrt{-s+4} - \sqrt{-s}}{\sqrt{-s+4} + \sqrt{-s}} \\s &= -\frac{(1-x)^2}{x}, \quad 0 \leq x \leq 1\end{aligned}$$

allows to change DEqs to the equations with factorized monomial denominators (!)

An example



e.g.:

$$\begin{aligned}
 A &= \frac{8 + s^2 - 2t + s(-6 + t + \epsilon t)}{(-4 + s)(-4 + s + t)} \\
 &\longrightarrow \frac{1}{(1+x)^2} \frac{1}{(x+y)} \frac{1}{(1+xy)} \times \{\dots\}
 \end{aligned}$$

Functions which appear in calculations

Now we have:

$$1\text{LBox} = \frac{2xy}{(1-x^2)(-1+y)^2} \left[\frac{1}{\epsilon} A + B + \epsilon C + \dots \right]$$

$$A = H[0, x]$$

$$B = H[0, x](H[0, y] + 2H[1, y])$$

$$C = f \{ \zeta_2, \zeta_3,$$

$$H[-1, x], H[0, x], H[-1, 0, x], H[-1, 0, 0, x], \dots,$$

$$H[0, y], H[1, y], H[0, 0, 0, y], \dots,$$

$$G[-y, x], G[-1/y, 0, x], G[-1/y, -1, 0, x]$$

$$G[-1/y, 0, 0, x], G[-y, -1, 0, x], G[-y, 0, 0, x], \dots, \}$$

If we can find solutions for ϵ^n , then it is easy to find solutions for $\epsilon^{(n+1)}$

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Vermaseren, Remiddi

$$H(0; x) = \ln x ,$$

$$H(1; x) = \int_0^x \frac{dx'}{1-x'} = -\ln(1-x) ,$$

$$H(-1; x) = \int_0^x \frac{dx'}{1+x'} = \ln(1+x)$$

$$H(\vec{0}_w; x) = \frac{1}{w!} \ln^w x$$

$$H(\vec{m}_w; x) = \int_0^x dx' f(a; x') H(\vec{m}_{w-1}; x')$$

Bhabha process:

$$f(0, x) = \frac{1}{x}, \quad f(1, x) = \frac{1}{1-x}, \quad f(-1, x) = \frac{1}{1+x}$$

Vermaseren, Remiddi

$$H(0; x) = \ln x, H(0, 0; x) = \int \frac{H[0, x]}{x} dx = \frac{1}{2!} \ln^2 x, \dots$$

$$H(1; x) = \int_0^x \frac{dx'}{1-x'} = -\ln(1-x),$$

$$H(-1; x) = \int_0^x \frac{dx'}{1+x'} = \ln(1+x)$$

$$H(0, 1; x) = \int_0^x \frac{dx'}{x'} H(1; x') = - \int_0^x \frac{dx'}{x'} \ln(1-x'),$$

$$H(0, 1; x) = \text{Li}_2(x),$$

$$H(1, 0; x) = \int_0^x \frac{dx'}{1-x'} H(0; x') = \int_0^x \frac{dx'}{1-x'} \ln x'$$

Algebra, IBP identities

$$\begin{aligned} H(1, 0; x) &= \int_0^x \frac{dx'}{1-x'} H(0; x') = \int_0^x \frac{dx'}{1-x'} \ln x' , \\ H(1, 0; x) &= -\ln x \ln(1-x) + \text{Li}_2(x) \\ H(1, 0; x) &= H(0, x)H(1, x) - H(0, 1; x) \end{aligned}$$

$$\begin{aligned} H(\vec{a}, x)H(\vec{b}, x) &= \sum_{\vec{c}=\vec{a}+\vec{b}} H(\vec{c}, x) \\ H(m_1, \dots, m_q, x) &= H(m_1, x)H(m_2, \dots, m_q, x) \\ &\quad - H(m_1, m_2, x)H(m_3, \dots, m_q, x) \\ &\quad + \dots + (-1)^{q+1}H(m_q, \dots, m_1, x) \end{aligned}$$

Not all functions can be found analytically

M. Czakon, JG, T. Riemann, PRD71,073009: four not solved functions of weight 4 (see also Vermaseren, Remiddi):

$$H(-1, 0, 0, 1, x) = I_1(x) = \int_0^x \frac{dy}{1+y} \text{Li}_3(y),$$

$$I_2(x) = \int_0^x \frac{dy}{1+y} \text{Li}_2(y) \ln(1-y),$$

$$I_3(x) = \int_0^x \frac{dy}{1+y} \ln(y) \ln^2(1-y),$$

$$I_4(x) = \int_0^x \frac{dy}{1+y} \ln^2(y) \ln(1-y).$$

Additional basis functions

Bonciani, Aglietti: 2 and 3 massive propagators for two loop electroweak vertices

Thresholds: $x = 0, x = 1, x = 4, \frac{1}{1-x^2+1}$

$$f(\mp 4; x) = \frac{1}{4 \pm x}, \quad f(c(\bar{c}), x) = \frac{1}{x - \frac{1}{2} \mp i \frac{\sqrt{3}}{2}}$$

$$f(\mp r, x) = \frac{1}{\sqrt{x(4 \pm x)}}, \quad f(\mp 1 \mp r, x) = \frac{1}{\sqrt{x(4 \pm x)} (1 \pm x)},$$

$$f(\pm 1/4; x) = \frac{1}{\frac{1}{4} \mp x}, \quad f(\pm 1 \pm r/4; x) = \frac{1}{\sqrt{x \mp \frac{1}{4}(1 \mp x)}},$$

$$f(r_0/4; x) = \frac{1 - 2i\sqrt{x - \frac{1}{4}}}{x\sqrt{x - \frac{1}{4}}}, \quad f(-r_0/4; x) = \frac{1 - 2\sqrt{x + \frac{1}{4}}}{x\sqrt{x + \frac{1}{4}}}$$

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$$g(-y, x) = \frac{1}{x+y}, \quad g(-1/y, x) = \frac{1}{x+\frac{1}{y}}$$

$$G(\{-y, -1/y\}, \dots, x) = \int_0^x dx' g(\{-y, -1/y\}, x') H(\dots, x')$$

e.g.:

$$G(-y; x) = \int_0^x \frac{dx'}{x' + y} = \ln\left(1 + \frac{x}{y}\right),$$

$$G(-1/y; x) = \int_0^x \frac{dx'}{x' + \frac{1}{y}} = \ln(1 + xy)$$

$$G(-y, 0; x) = \int_0^x \frac{dx'}{x' + y} H(0, x')$$

Additional basis functions: Glover et al., massless two-loop diagrams with 3 off-shell legs

Kinematics:

$$\begin{aligned}\lambda &= \sqrt{(1-x-y)^2 - 4xy} = \sqrt{(x-x_0)(x-x_1)} \\ x_0 &= (1-\sqrt{y})^2, \quad x_1 = (1+\sqrt{y})^2\end{aligned}$$

$$f(\lambda, x) = \frac{1}{\lambda}, \quad f(x\lambda, x) = \frac{1}{x\lambda},$$

$$f(x_0, x) = -\frac{1}{x-x_0}, \quad f(x_1, x) = -\frac{1}{x-x_1}$$

$$\text{e.g. : } G(0; x, y) = \log \left(\frac{x}{x_0} \right),$$

$$G(\lambda; x, y) = \log \left(\frac{1-x+y-\lambda}{2\sqrt{y}} \right)$$

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GPLs and HPLs to the weight 4, set={-1,0,1,-y,-1/y}

weight 1, #5	weight 2, #25	weight 3, #125	weight 4, #125
2 GPLs	16 GPLs	98 GPLs	544 GPLs
3 HPLs	9 HPLs	27 HPLs	81 HPLs

```
/*
weight 2, 16 GPLs, 9 HPLs = 25

irreducible: G[-y_,-1,x_],
G[-y_,-1,x_],
G[-y_,0,x_],
G[-y_,1,x_],
G[-y_,-1/y,x_]

*)
```

$$G[-y_,-1,x_]:= (\text{Log}[1+x]*\text{Log}[(x+y)/(-1+y)] - \text{PolyLog}[2,(1-y)^{(-1)}] + \text{PolyLog}[2,(1+x)/(1-y)])$$

$$G[-1/y_,-1,x_]:= G[-y,-1,x]/. y \rightarrow 1/y$$

$$G[-y_,-y_,-1]:= 1/2*G[-y,x]^2$$

$$\dots$$

$$G[-y_,-1,x_]:= ((\text{Log}[1-x]^2*\text{Log}[(x+y)/(1+y)])/2 + \text{Log}[1-x]*\text{PolyLog}[2,(1-x)/(1+y)] + \text{PolyLog}[3,(1+y)^{(-1)}] - \text{PolyLog}[3,(1-x)/(1+y)])$$

$$G[-y_,-(y_),-1,x_]:= \text{NIntegrate}[G[-y,-1,x]/(y+z),\{z,0,x\}]$$

$$\dots$$

$$G[-y_,-(y_),-(y_),-1,x_]:= \text{NIntegrate}[\text{If}[x2 < x1, G[-1,-1,x2] v^2/((1+v*x1)*(1+v*x2)) /. v \rightarrow y, 0], \{x1,0,x\}, \{x2,0,x1\}]$$

$$\dots$$

```
(*  

weight 2, 16 GPLs, 9 HPLs = 25

irreducible: G[-y_, x_],  

G[-y_, -1,x_],  

G[-y_, 0,x_],  

G[-y_, 1,x_],  

G[-y_, -1/y,x_]  

*)  

G[-y_,-1,x_]:= (Log[1+x]*Log[(x+y)/(-1+y)] -  

PolyLog[2, (1-y)^(-1)] +  

PolyLog[2, (1+x)/(1-y)] )  

G[-1/y_,-1,x_]:= G[-y,-1,x] /. y -> 1/y  

G[-y_,-y_,x_]:= 1/2*G[-y,x]^2  

...  

G[-y_, 1, 1,x_]:= ((Log[1-x]^2*Log[(x+y)/(1+y)])/2 +  

Log[1-x]*PolyLog[2, (1-x)/(1+y)] +  

PolyLog[3, (1+y)^(-1)] -  

PolyLog[3, (1-x)/(1+y)])  

G[-y_, -(y_), -1,x_]:= NIntegrate[G[-y,-1,z]/(y+z),{z,0,x}];  

...  

G[-(y_)^(-1),-(y_)^(-1),-1,-1,x_]:=  

NIntegrate[ If[z2 < z1, G[-1,-1,z2]v^2/((1+v*z1)*(1+v*z2))  

/.v -> y, 0], {z1,0,x},{z2,0,z1}];  

...  


```

► Relations between GPLs in s and t channels

$$\begin{aligned} G(-y, -1, x) \longrightarrow & H(-1, x)H(1, y) - H(0, x)H(1, y) \\ & + H(-1, x)H(-x, y) + H(0, -1, x) \\ & + H(1, 0, y) - G(1, -x, y) \end{aligned}$$

► Two different transformations

$$x = \frac{\sqrt{-s+4} - \sqrt{-s}}{\sqrt{-s+4} + \sqrt{-s}} \quad \text{and} \quad z = \frac{1}{\sqrt{1 - 4/s}}$$

then

$$x = \frac{1-z}{1+z}$$



$$x = \frac{1-x'}{1+x'}, \quad y = \frac{1-y'}{1+y'}$$

$$H\left(-1, -1, \frac{1-x}{1+x}\right) \longrightarrow H(-1, -1, x')$$

$$-H(-1, x')H(-1, 1) + H(-1, -1, 1)$$

$$G(-y, 0, x) \longrightarrow H\left(-\frac{1-y'}{1+y'}, 0, 1\right) + \int_1^x dz \frac{H(0, z)}{x' + \frac{1-y'}{1+y'}}$$

$$= \{-\zeta_2 + H(0, -1, y') - H(0, 0, y')\}$$

$$+ \int_0^{x'} dz [H(1, z) + H(-1, z)] \times \left[\frac{1}{1+z} + \frac{y'}{1-zy'} \right]$$

$$= \{\dots\} + \dots + H(-1, 1, x') + G(1/y', -1, x) + G(1/y', 1, x)$$

- ▶ Formally, generating functions with indices {-1,0,1,-y,-1/y} must be extended by generating functions with indices +y, +1/y

- ▶ Calculations of Feynman integrals with DEqs method:
new functions
- ▶ results can be represented analytically in an easy way in
order by order ϵ parameter expansion
- ▶ Mathematica codes for algebraic manipulations [$x \leftrightarrow y$,
[conformal transformations](#)] and [numerical calculation](#) of
HPLs and GPLs are in use for Bhabha calculations
(Fortran code by Gehrmann and Remiddi for numerical
calculation of HPLs and GPLs exists either)